NON-NEGATIVE PERTURBATIONS OF NON-NEGATIVE SELF-ADJOINT OPERATORS

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In memory of my great teacher Mark Krein.

ABSTRACT. Let A be a non-negative self-adjoint operator in a Hilbert space \mathcal{H} and A_0 be some densely defined closed restriction of A_0 , $A_0 \subseteq A \neq A_0$. It is of interest to know whether A is the unique non-negative self-adjoint extensions of A_0 in \mathcal{H} . We give a natural criterion that this is the case and if it fails, we describe all non-negative extensions of A_0 . The obtained results are applied to investigation of non-negative singular point perturbations of the Laplace and poly-harmonic operators in $\mathbb{L}_2(\mathbf{R}_n)$.

1. Introduction

In this paper we deal with a non-negative self-adjoint operator A in a Hilbert space \mathcal{H} , some densely defined not essentially self-adjoint restriction A_0 of A and again with self-adjoint extensions of A_0 in \mathcal{H} , which following [1] we call here *singular perturbations* of A. For quick getting onto the matter of the main problem let us compare the point perturbations of self-adjoint Laplace operators $-\Delta$ in three and two dimensions acting in $\mathbb{L}_2(\mathbf{R}_3)$ and $\mathbb{L}_2(\mathbf{R}_2)$, respectively, that is let us consider the restriction $-\Delta^0$ of $-\Delta$ onto the Sobolev subspaces $\mathbb{H}_2^2(\mathbf{R}_i \setminus \{0\})$, i = 3, 2, and self-adjoint extensions $-\Delta_{\alpha}$, $\alpha \in \mathbf{R}$, of $-\Delta^0$ in $\mathbb{L}_2(\mathbf{R}_i)$ with domains

$$\mathcal{D}_{\alpha}^{(3)} := \left\{ f : f \in \mathbb{H}_{2}^{2}(\mathbf{R}_{3}), \lim_{|\mathbf{x}| \downarrow 0} \left[\frac{d}{d|\mathbf{x}|} (|\mathbf{x}| f(\mathbf{x})) - \alpha |\mathbf{x}| f(\mathbf{x}) \right] = 0 \right\},$$

$$\mathcal{D}_{\alpha}^{(2)} := \left\{ f : f \in \mathbb{H}_{2}^{2}(\mathbf{R}_{2}), \lim_{|\mathbf{x}| \downarrow 0} \left[\left(\frac{2\pi\alpha}{\ln |\mathbf{x}|} + 1 \right) f(\mathbf{x}) - \lim_{|\mathbf{x}'| \downarrow 0} \frac{\ln |\mathbf{x}|}{\ln |\mathbf{x}'|} f(\mathbf{x}') \right] = 0 \right\}.$$

The self-adjoint operators $-\Delta_{\alpha}$ are just mentioned above singular perturbations of $-\Delta$. The resolvents $(-\Delta_{\alpha}-z)^{-1}$, $z\in\rho(-\Delta_{\alpha})$, of the operators $-\Delta_{\alpha}$ act in the corresponding spaces \mathbb{L}_2 as integral operators with kernels (Green functions) [1],

$$(1.2) G_{\alpha,z}^{3}(\mathbf{x},\mathbf{x}') = \begin{cases} G_{z}^{(0)}(\mathbf{x},\mathbf{x}') + (\alpha - i\sqrt{z}/4\pi)^{-1}G_{z}^{(0)}(\mathbf{x},0)G_{z}^{(0)}(0,\mathbf{x}'), \\ G_{z}^{(0)}(\mathbf{x},\mathbf{x}') = \frac{\exp i\sqrt{z}|\mathbf{x}-\mathbf{x}'|}{4\pi|\mathbf{x}-\mathbf{x}'|} \text{ (three dimension)}; \end{cases}$$

(1.3)
$$G_{\alpha,z}^{2}(\mathbf{x}, \mathbf{x}') = \begin{cases} G_{z}^{(0)}(\mathbf{x}, \mathbf{x}') + 2\pi(2\pi\alpha - \psi(1) + \ln\sqrt{z}/2i)^{-1}G_{z}^{(0)}(\mathbf{x}, 0)G_{z}^{(0)}(0, \mathbf{x}'), \\ G_{z}^{(0)}(\mathbf{x}, \mathbf{x}') = (\frac{i}{4})H_{0}^{(1)}(i\sqrt{z}|\mathbf{x} - \mathbf{x}'|) \text{ (two dimension)}. \end{cases}$$

By (1.2) the Green function $G_{\alpha,z}(\mathbf{x},\mathbf{x}')$ of the self-adjoint operator $-\Delta_{\alpha}$ in $\mathbb{L}_2(\mathbf{R}_3)$ is holomorphic on the half-axis $(-\infty,0)$ for $\alpha \geq 0$ and has on this half-axis a simple pole

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for $\alpha < 0$. Hence in the case of three dimensions self-adjoint extensions $-\Delta_{\alpha}$ are non-negative for all $(\alpha \geq 0)$ and non-positive for $\alpha < 0$.

Contrary to this by (1.3) in the case of two dimensions for any $\alpha \in \mathbf{R}$ the Green function $G_{\alpha,z}$ has a simple pole on the half-axis $(-\infty,0)$. Hence all singular perturbations $-\Delta_{\alpha}$ of the two-dimensional Laplace operators have one negative eigenvalue. In other words the Laplace operator $-\Delta$ defined in a standard way is a unique non-negative self-adjoint extension in $\mathbb{L}_2(\mathbf{R}_2)$ of the symmetric operator $-\Delta^0$ in $\mathbb{L}_2(\mathbf{R}_2)$.

In this note we try to reveal the underlying cause of such discrepancy. Recall that each densely defined non-negative symmetric operator has at least one non-negative canonical self-adjoint extension (Friedrichs extension). In a more general setting we try to understand here why in some cases the non-negative extension appears to be unique. Actually this questions is embedded into the framework of the general extension theory for semi-bounded symmetric operators developed in the famous paper of M. G. Krein [2]. Naturally, there is a criterion of uniqueness of non-negative extension in [2]. In the next Section, using only approaches of [2] we find another form of this criterion, directly facilitating the investigation of singular perturbations and, for cases where the conditions of these criteria fail, describe all non-negative singular perturbations of a given non-negative self-adjoint operator A associated with some its densely defined non-self-adjoint restriction A_0 . In fact we give here a parametrization of the operator interval $[A_{\mu}, A_M]$ of all canonical non-negative self-adjoint extensions of a given densely defined non-negative operator. The third Section illustrates the obtained results by an example of singular perturbations of the Laplace and poly-harmonic operators in $\mathbb{L}_2(\mathbf{R}_n)$.

Note that very close results were obtained recently in a somewhat different way in [3] where, in terms of this note, the authors have described singular perturbations of the Friedrichs extension of a given densely defined non-negative operator and also gave an illustration with an example of singular perturbations of the Laplace operator in $\mathbb{L}_2(\mathbf{R}_3)$.

2. Uniqueness criterion and parametrization of non-negative singular perturbations

Let A be a non-negative self-adjoint operator acting in the Hilbert space \mathcal{H} and A_0 be a densely defined closed operator, which is a restriction of A onto a subset $\mathcal{D}(A_0)$ of the domain $\mathcal{D}(A)$ of A. Let us consider the subspaces $\mathcal{M} := (I + A_0 \mathcal{D}(A_0))$ and $\mathcal{N} := \mathcal{H} \ominus \mathcal{M}$. We will assume that

$$(2.1) 1) \mathcal{M} \neq \mathcal{H}, 2) \mathcal{N} \cap \mathcal{D}(A) = \{0\}.$$

We call all self-adjoint extensions of A_0 in \mathcal{H} other than the given A singular perturbations of A. It is of interest to know whether there are non-negative operators among singular perturbations of A. In this section we try to find a convenient criterion that such singular perturbations of A do not exist. In other words we look for a criterion that A is the only one non-negative operator among all self-adjoint extensions of A_0 . Following the approach developed in the renowned paper of M. G. Krein [2] let us consider the operator from $K_0: \mathcal{M} \to \mathcal{H}$ defined by the relations

(2.2)
$$f = (I + A_0) x, \quad K_0 f = A_0 x, \quad x \in \mathcal{D}(A_0).$$

It is easy to see that K_0 is a non-negative contraction.

$$(2.3) (K_0 f, f) \ge 0, ||K_0 f||^2 \le ||f||^2, f \in \mathcal{M}.$$

Let A_1 be any non-negative self-adjoint extension of A_0 in \mathcal{H} . Then $K_1 := A_1 (A_1 + I)^{-1}$ is a non-negative operator, which is a contractive extension of K_0 from the domain \mathcal{M} onto the whole \mathcal{H} , $K_1f = K_0f$, $f \in \mathcal{M}$.

¹The attention of the author to this phenomenon was drawn by Sergey Gredeskul.

On the other hand, for any contractive extension K_1 from \mathcal{M} onto \mathcal{H} such that the unity is not its eigenvalue the non-negative self-adjoint operator $A_1 = K_1 (I - K_1)^{-1}$ is a self-adjoint extension of A_0 in \mathcal{H} . Therefore, A_0 has a unique non-negative self-adjoint extension in \mathcal{H} if and only if K_0 admits only one non-negative contractive extension onto the whole \mathcal{H} , no eigenvalue of which equals 1, that is, $K = A(I + A)^{-1}$. So the uniqueness of A as a non-negative extension of A_0 is equivalent to uniqueness of K_0 as a non-negative contractive extension of K_0 .

From now on we will denote by \mathbf{G} the set consisting of A and all its singular perturbations and by \mathbf{C} the set of non-negative contractions obtained from \mathbf{G} by the transformation $A_1 \to A_1 (A_1 + I)^{-1}$, $A_1 \in \mathbf{G}$. Let us denote by $P_{\mathcal{M}}$ the orthogonal projector onto \mathcal{M} in \mathcal{H} and let $P_{\mathcal{N}} = I - P_{\mathcal{M}}$. With respect to representation of \mathcal{H} as the orthogonal sum $\mathcal{M} \oplus \mathcal{N}$ we can represent each operator from \mathbf{C} as 2×2 block operator matrix

(2.4)
$$K_X = \begin{pmatrix} T & \Gamma^* \\ \Gamma & X \end{pmatrix}.$$

Here

$$T = P_{\mathcal{M}} K_0|_{\mathcal{M}}, \quad \Gamma = P_{\mathcal{M}} K_0|_{\mathcal{M}}$$

and X is some non-negative contraction in \mathcal{N} , which distinguishes different elements from C. Since each $K_X \in \mathbf{C}$ is non-negative and contractive, we have

$$(2.5) T \ge 0, \quad T^2 + \Gamma^* \Gamma \le I.$$

Note further that $K_X \in \mathbf{C}$ is equivalent to

$$(2.6) K_X + \varepsilon I \ge 0, \quad (1 + \varepsilon)I - K_X \ge 0$$

for any $\varepsilon > 0$.

The block matrix representation of K_X and the Schur-Frobenius factorization formula transform (2.6) into the following block matrix inequalities:

(2.7)
$$\begin{pmatrix} I & 0 \\ \Gamma(T+\varepsilon)^{-1} & I \end{pmatrix} \begin{pmatrix} T+\varepsilon & 0 \\ 0 & X+\varepsilon-\Gamma(T+\varepsilon)^{-1}\Gamma^* \\ \times \begin{pmatrix} I & (T+\varepsilon)^{-1}\Gamma^* \\ 0 & I \end{pmatrix} \geq 0,$$

$$(2.8) \qquad \begin{pmatrix} I & 0 \\ -\Gamma(I+\varepsilon-T)^{-1} & I \end{pmatrix} \begin{pmatrix} 1+\varepsilon-T & 0 \\ 0 & 1+\varepsilon-X-\Gamma(1+\varepsilon-T)^{-1}\Gamma^* \\ \times \begin{pmatrix} I & -(1+\varepsilon-T)^{-1}\Gamma^* \\ 0 & I \end{pmatrix} \ge 0.$$

By our assumptions $T \ge 0$ and $I - T \ge 0$. Therefore block matrix inequalities (2.7) and (2.8) are reduced to

$$\begin{cases} X + \varepsilon I - \Gamma(T + \varepsilon I)^{-1} \Gamma^* \ge 0, \\ (1 + \varepsilon)I - X - \Gamma[(1 + \varepsilon)I - T]^{-1} \Gamma^* \ge 0, \ \varepsilon > 0. \end{cases}$$

Observe that the operator functions of ε in the left-hand sides of inequalities (2.9) are monotone. Setting

$$Y := X - \lim_{\varepsilon \downarrow 0} \Gamma(T + \varepsilon I)^{-1} \Gamma^*$$

we conclude from (2.9) that $K_X \in \mathbf{C}$ if and only if

(2.10)
$$0 \le Y \le I - \lim_{\varepsilon \downarrow 0} \left\{ \Gamma(T + \varepsilon I)^{-1} \Gamma^* + \Gamma[(1 + \varepsilon)I - T]^{-1} \Gamma^* \right\}.$$

Hence the equality

(2.11)
$$I - \lim_{\varepsilon \downarrow 0} \left\{ \Gamma(T + \varepsilon I)^{-1} \Gamma^* + \Gamma[(1 + \varepsilon)I - T]^{-1} \Gamma^* \right\} = 0$$

is a criterion that there are no contractive non-negative extensions of K_0 in \mathcal{H} other than K.

Let us express now (2.10) and (2.11) in terms of given K and A. To this end we use the following proposition.

Proposition 2.1. Let L be a bounded invertible operator in the Hilbert space $\mathcal{H} = \mathcal{M} \oplus \mathcal{N}$ given as $2 \times @$ block operator matrix

$$L = \left(\begin{array}{cc} R & U \\ V & S \end{array} \right),$$

where R and S are invertible operators in \mathcal{M} and \mathcal{N} , respectively, and U, V act between \mathcal{M} and \mathcal{N} . If R is an invertible operator in \mathcal{M} , then

(2.12)
$$\begin{pmatrix} R^{-1} & 0 \\ 0 & 0 \end{pmatrix} = L^{-1} - L^{-1} P_{\mathcal{N}} \Lambda^{-1} P_{\mathcal{N}} L^{-1}, \quad \Lambda = P_{\mathcal{N}} L^{-1}|_{\mathcal{N}}.$$

Setting

(2.13)
$$\Lambda_{1,\varepsilon} = P_{\mathcal{N}}(K+\varepsilon I)^{-1}|_{\mathcal{N}}, \quad \Lambda_{2,\varepsilon} = P_{\mathcal{N}}[(1+\varepsilon)I - K]^{-1}|_{\mathcal{N}}$$

and applying (2.12) with $L = K + \varepsilon I$ and

$$\begin{split} R &= T + \varepsilon I, & U &= \Gamma^* = P_{\mathcal{M}} K|_{\mathcal{N}} = P_{\mathcal{M}}[K + \varepsilon I]|_{\mathcal{N}}, \\ V &= \Gamma = P_{\mathcal{N}} K|_{\mathcal{M}} = P_{\mathcal{N}}[K + \varepsilon I]|_{\mathcal{M}}, & S &= P_{\mathcal{N}} K|_{\mathcal{N}} + \varepsilon I, \end{split}$$

yields

$$\Gamma(T+\varepsilon I)^{-1}\Gamma^* = P_{\mathcal{N}}K|_{\mathcal{N}} + \varepsilon I - \Lambda_{1,\varepsilon}^{-1}$$

In the same fashion we get

$$\Gamma[(1+\varepsilon)I - T]^{-1}\Gamma^* = P_{\mathcal{N}}[I - K]|_{\mathcal{N}} + \varepsilon I - \Lambda_{2,\varepsilon}^{-1}.$$

Hence

$$(2.14) I - \lim_{\varepsilon \downarrow 0} \left(\Gamma(T + \varepsilon I)^{-1} \Gamma^* + \Gamma[(1 + \varepsilon)I - T]^{-1} \Gamma^* \right) = \lim_{\varepsilon \downarrow 0} \Lambda_{1,\varepsilon}^{-1} + \lim_{\varepsilon \downarrow 0} \Lambda_{2,\varepsilon}^{-1}.$$

Combining (2.10), (2.11) and (2.14) results in the following theorem.

Theorem 2.2. Let K be a non-negative contraction in the Hilbert space $\mathcal{H} = \mathcal{M} \oplus \mathcal{N}$, K_0 is the restriction of K onto the subspace $\mathcal{M}(=\mathcal{M} \oplus \{0\})$ and

$$G_1 = \lim_{\varepsilon \downarrow 0} (P_{\mathcal{N}}[K + \varepsilon I]|_{\mathcal{N}})^{-1}, \quad G_2 = \lim_{\varepsilon \downarrow 0} (P_{\mathcal{N}}[I - K + \varepsilon I]|_{\mathcal{N}})^{-1}.$$

Then the set C of all non-negative contractive extensions K_X of K_0 in \mathcal{H} is described by the expression

(2.15)
$$K_X = \begin{pmatrix} P_{\mathcal{M}} K|_{\mathcal{M}} & P_{\mathcal{M}} K|_{\mathcal{N}} \\ P_{\mathcal{M}} K|_{\mathcal{N}} & X \end{pmatrix},$$

where X runs over the set of all non-negative contractions in $\mathcal N$ satisfying the inequalities

$$(2.16) P_{\mathcal{N}}K|_{\mathcal{N}} - G_1 \le X \le P_{\mathcal{N}}K|_{\mathcal{N}} + G_2.$$

In particular, K is a unique non-negative contractive extension of K_0 if and only if $G_1 = G_2 = 0$.

Remark 2.3. The set **C** of non-negative contractions of K_0 contains a minimal extension $K_{X_{\mu}}$ with $X_{\mu} = P_{\mathcal{N}}K|_{\mathcal{N}} - G_1$ in (2.15) and a maximal extension K_{X_M} with $X_M = P_{\mathcal{N}}K|_{\mathcal{N}} + G_2$ in (2.15). If $G_1 = 0$ ($G_2 = 0$), then K is a minimal (maximal) element of \mathbf{C}

Theorem 2.2 can be formulated in terms of a non-negative self-adjoint operator A and its non-negative singular perturbations.

Theorem 2.4. Let A be a non-negative self-adjoint operator in a Hilbert space \mathcal{H} , A_0 is a densely defined closed symmetric operator, which is a restriction of A onto a linear subset $\mathcal{D}(A_0) \subset \mathcal{D}(A)$ such that $\mathcal{N} = (I + A)\mathcal{D}(A_0) \neq \{0\}$ and let

$$G_1 = \lim_{\varepsilon \downarrow 0} \left(P_{\mathcal{N}}[I+A][A+\varepsilon I]|_{\mathcal{N}} \right)^{-1}, \quad G_2 = \lim_{\varepsilon \downarrow 0} \left(P_{\mathcal{N}}[I+A][I+\varepsilon A]|_{\mathcal{N}} \right)^{-1}.$$

Then the set of all non-negative singular perturbations A_Y of A is described by the formula

(2.17)
$$\begin{cases} f = g - Y(I+A)g, \\ A_Y f = Ag + Y(I+A)g, \end{cases}$$

where $g \in \mathcal{D}(A)$ and Y runs over the set of non-negative contractions in \mathcal{N} satisfying the inequalities

$$(2.18) -G_1 \le Y \le G_2.$$

A has no singular non-negative perturbations if and only if $G_1 = G_2 = 0$.

Remark 2.5. The set of all non-negative singular perturbations of A contains a minimal perturbation A_{μ} with and a maximal perturbation A_{M} such that any non-negative perturbation A_{X} satisfies the inequalities

$$(I + A_M)^{-1} \le A_Y \le (I + A_\mu)^{-1}$$
.

The corresponding values of the parameters Y in Theorem 2.4 are

(2.19)
$$Y_{\mu} = -G_1, Y_M = G_2.$$

If $G_1 = 0$ ($G_2 = 0$), then the minimal (maximal) perturbation coincides with A.

By simple calculation we get from (2.17) the following version of the M. G. Krein resolvent formula.

Proposition 2.6. The set of resolvents of all non-negative singular perturbations A_Y of A is described by the M. G. Krein formula (2.20)

$$(A_Y - zI)^{-1} = (A - zI)^{-1}$$
$$- (1 + z)(A + I)(A - zI)^{-1}Y \left[I + (1 + z)P_{\mathcal{N}}(A + I)(A - zI)^{-1}Y \right]^{-1}$$
$$\times P_{\mathcal{N}}(A + I)(A - zI)^{-1},$$

where Y runs over the contractions in \mathcal{N} satisfying inequalities $-G_1 \leq Y \leq G_2$.

3. Application to some differential operators

Let us consider the multiplication operator A in $\mathbb{L}_2(\mathbf{R}_n)$ by a continuous function $\varphi(k)$, $k^2 = k_1^2 + \cdots + k_n^2$, such that $\varphi(k) > 0$ almost everywhere and

$$\int_0^\infty \frac{1}{(1+\varphi(k))^2} k^{n-1} dk < \infty.$$

A is a non-negative self-adjoint operator,

$$\mathcal{D}(A) = \left\{ f : \int_{\mathbb{R}} |1 + \varphi(k)|^2 |f(\mathbf{k})|^2 d\mathbf{k} < \infty, \ f \in \mathbb{L}_2(\mathbf{R}_n) \right\}.$$

In the sequel $\hat{\delta}$ stands for the unbounded linear functional in $\mathbb{L}_2(\mathbf{R}_n)$, formally defined as follows:

$$\hat{\delta}(f) = \int_{\mathbb{R}_n} f(\mathbf{k}) d\mathbf{k}.$$

Note that the domain of $\hat{\delta}$ contains $\mathcal{D}(A)$. Let us denote by A_0 the restriction of A onto the linear set

(3.2)
$$\mathcal{D}_0(A) := \left\{ f : f \in \mathcal{D}(A), \ \hat{\delta}(f) = 0 \right\}.$$

The closure of $A_0 \neq A$ and

$$\mathcal{N} = (\mathbb{L}_2(\mathbf{R}_n) \ominus (I+A)\mathcal{D}_0(A)) = \left\{ \xi \cdot \frac{1}{1+\varphi(k)}, \ \xi \in \mathbf{C} \right\}.$$

Applying Theorem 2.4 yields

Proposition 3.1. A is a unique non-negative self-adjoint extension of A_0 , that is, A has no non-negative singular perturbations if and only if

(3.3)
$$\int_0^\infty \frac{1}{\varphi(k)(1+\varphi(k))} k^{n-1} dk = \infty \quad \text{and} \quad \int_0^\infty \frac{1}{(1+\varphi(k))} k^{n-1} dk = \infty.$$

Put $\varphi(k) = k^2$ and let n = 2. Then the both integrals in Proposition 3.1 are divergent. Hence the restriction A_0 of the operator A of multiplication by k^2 in $\mathbb{L}_2(\mathbf{R}_2)$ onto the linear set (3.2) has a unique non-negative self-adjoint extension in \mathbb{L}_2 . Note that the multiplication operator by k^2 in $\mathbb{L}_2(\mathbf{R}_n)$ is isomorphic to the self-adjoint Laplace operator $-\Delta$ in $\mathbb{L}_2(\mathbf{R}_n)$ and its considered here restriction A_0 is isomorphic to the restriction of $-\Delta$ onto the Sobolev subspace $\mathbb{H}_2^2(\mathbf{R}_n \setminus \{0\})$. It follows that the self-adjoint Laplace operator in $\mathbb{L}_2(\mathbf{R}_2)$ has no non-negative singular perturbations with support at one point of \mathbf{R}_2 .

However, the non-negative singular perturbations of $-\Delta$ in $\mathbb{L}_2(\mathbf{R}_2)$ with support at two or more points do already exist. For example, let us consider there the restriction A_0 of the multiplication operator by k^2 , for which the defect subspace \mathcal{N} is one-dimensional and consists of functions collinear to

$$e_0(\mathbf{k}) = \frac{1 - \exp(-i(\mathbf{k} \cdot \mathbf{x}_0))}{1 + k^2}, \quad \mathbf{x}_0 \in \mathbf{R}_2.$$

In this case,

$$||e_{0}||^{2} = \int_{\mathbf{R}_{2}} \frac{4\sin^{2}\frac{1}{2}(\mathbf{k} \cdot \mathbf{x}_{0})}{(1+k^{2})^{2}} \cdot d\mathbf{k} < \infty,$$

$$((I+A)A^{-1}e_{0}, e_{0}) = \int_{\mathbf{R}_{2}} \frac{4\sin^{2}\frac{1}{2}(\mathbf{k} \cdot \mathbf{x}_{0})}{k^{2}(1+k^{2})} \cdot d\mathbf{k} < \infty,$$

$$((I+A)e_{0}, e_{0}) = \int_{\mathbf{R}_{2}} \frac{4\sin^{2}\frac{1}{2}(\mathbf{k} \cdot \mathbf{x}_{0})}{1+k^{2}} \cdot d\mathbf{k} = \infty.$$

Hence $G_1 = ||e_0||^2 \cdot ((I+A)e_0, e_0)^{-1} > 0$, but $G_2 = 0$. Hence, the considered restriction A_0 of the multiplication operator A by k^2 has non-negative self-adjoint extensions in $\mathbb{L}_2(\mathbf{R}_2)$ other than A, and A is a maximal element in the set of these extensions. It remains to note that A is isomorphic to the self-adjoint Laplace operator $-\Delta$ in $\mathbb{L}_2(\mathbf{R}_2)$, and A_0 is isomorphic to the restriction of this $-\Delta$ to the subset of functions $f(\mathbf{x})$ from $\mathcal{D}(-\Delta)$ satisfying the conditions

$$\begin{split} &\lim_{|\mathbf{x}| \to 0} (\ln |\mathbf{x}|)^{-1} f(\mathbf{x}) - \lim_{|\mathbf{x} - \mathbf{x}_0| \to 0} (\ln |\mathbf{x} - \mathbf{x}_0|)^{-1} f(\mathbf{x}) = 0, \\ &\lim_{|\mathbf{x}| \to 0} \left[f(\mathbf{x}) - \ln |\mathbf{x}| \lim_{|\mathbf{x}'| \to 0} (\ln |\mathbf{x}'|)^{-1} f(\mathbf{x}') \right] \\ &- \lim_{|\mathbf{x} - \mathbf{x}_0| \to 0} \left[f(\mathbf{x}) - \ln |\mathbf{x} - \mathbf{x}_0| \lim_{|\mathbf{x}' - \mathbf{x}_0| \to 0} (\ln |\mathbf{x}' - \mathbf{x}_0|)^{-1} f(\mathbf{x}') \right] = 0. \end{split}$$

Put now as above $\varphi(k) = k^2$ and let n = 3. Then the first integral in Proposition 3.1 is convergent while the second one, as before, is divergent. Hence the restriction A_0 of the

operator A of multiplication by k^2 in $\mathbb{L}_2(\mathbf{R}_3)$ onto the linear set (3.2) has infinitely many non-negative self-adjoint extension in $\mathbb{L}_2(\mathbf{R}_3)$. Hence, the self-adjoint Laplace operator in $\mathbb{L}_2(\mathbf{R}_3)$ has infinitely many non-negative singular perturbations with support at one point of \mathbf{R}_3 and the standard Laplace operator is a maximal element in the set of these perturbations.

As the next example we consider the multiplication operator A by k^{2l} in $\mathbb{L}_2(\mathbf{R}_n)$ assuming that $4l \leq n+1$. A is isomorphic to the polyharmonic operator $(-\Delta)^l$ in $\mathbb{L}_2(\mathbf{R}_n)$. Let us consider the restriction A_0 of A with the domain (3.2) that is nonnegative symmetric operator which is isomorphic to the restriction of the polyharmonic operator $(-\Delta)^l$ onto the Sobolev subspace $\mathbb{H}^2_{2l}(\mathbf{R}_n \setminus \{0\})$. Applying Theorem 2.4 and Proposition 3.1 results in the following proposition.

Proposition 3.2. If n < 2l then there are infinitely many non-negative singular perturbations of $(-\Delta)^l$ associated with the one-point symmetric restriction A_0 and $(-\Delta)^l$ is a minimal element in the set of non-negative extensions of A_0 in $\mathbb{H}^2_{2l}(\mathbf{R}_n \setminus \{0\})$.

If n=2l then $(-\Delta)^l$ has no such perturbations in $\mathbb{H}^2_{2l}(\mathbf{R}_n\setminus\{0\})$.

If n > 2l then there is an infinite set of non-negative singular perturbations of $(-\Delta)^l$ associated with A_0 and, for them considered as non-negative extensions of A_0 in the set $\mathbb{H}^2_{2l}(\mathbf{R}_n \setminus \{0\})$, the operator $(-\Delta)^l$ is a maximal element.

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