

BOUNDARY PROBLEMS FOR FULLY NONLINEAR PARABOLIC EQUATIONS WITH LÉVY LAPLACIAN

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ABSTRACT. We suggest a method to solve boundary and initial-boundary value problems for a class of nonlinear parabolic equations with the infinite dimensional Lévy Laplacian Δ_L

$$f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0$$

in fundamental domains of a Hilbert space.

1. INTRODUCTION

In the paper by S. Albeverio, Ya. I. Belopolskaya, M. N. Feller [1] we have constructed a solution of the Cauchy problem for a fully nonlinear parabolic equation with the Lévy Laplacian

$$f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0, \quad U(0, x) = U_0(x),$$

where $f(\xi, \eta, \zeta)$ is a function on R^3 .

In the present paper we continue the investigation started in [1]. We develop a method to solve the boundary value problem

$$f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0, \quad U(t, x) = G(t, x) \quad \text{on } \Gamma,$$

and the initial-boundary value problem

$$f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0, \quad U(0, x) = U_0(x), \quad U(t, x) = G(t, x) \quad \text{on } \Gamma,$$

in fundamental domains $\Omega \cup \Gamma$ of a Hilbert space.

2. PRELIMINARIES

Let H be a separable real Hilbert space, $F(x)$ be a scalar function defined on H .

An infinite dimensional Laplacian was introduced by P. Lévy [2] through the formula

$$\Delta_L F(x) = 2 \lim_{\rho \rightarrow 0} \frac{\mathfrak{M}F(x + \rho y) - F(x)}{\rho^2},$$

where $\mathfrak{M}\Phi$ is the mean value of the function $\Phi(y)$ over the sphere $\|y\|_H^2 = 1$.

If $F(x)$ is a twice strongly differentiable function at the point x_0 , then the Lévy Laplacian is defined (when it exists) by the formula

$$(1) \quad \Delta_L F(x_0) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n (F''(x_0) f_k, f_k)_H,$$

where $F''(x)$ is the Hessian of the function $F(x)$, $\{f_k\}_1^\infty$ is an orthonormal basis in H .

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In the sequel we need a property of the Lévy Laplacian that has been established in [2], (see as well [3]) which we describe now.

Let

$$F(x) = f(U_1(x), \dots, U_m(x)),$$

$f(u_1, \dots, u_m)$ be a twice continuously differentiable function with m arguments defined on the domain $\{U_1(x), \dots, U_m(x)\} \subset R^m$, where $(U_1(x), \dots, U_m(x))$ is a vector of values of the functions $U_1(x), \dots, U_m(x)$. Assume that $U_j(x)$ are uniformly continuous in a bounded domain $\Omega \subset H$ and twice strongly differentiable functions and $\Delta_L U_j(x)$ exist ($j = 1, \dots, m$). Then $\Delta_L F(x)$ exists and

$$(2) \quad \Delta_L F(x) = \sum_{j=1}^m \frac{\partial f}{\partial u_j} \Big|_{u_j=U_j(x)} \Delta_L U_j(x).$$

Actually the second differential of the function $F(x)$ at the point x in the direction $h \in H$ has the form

$$\begin{aligned} d^2 F(x; h) &= (F''(x)h, h)_H = \sum_{i,j=1}^m \frac{\partial^2 f}{\partial u_i \partial u_j} \Big|_{u_l=U_l(x)} (U'_i(x), h)_H (U'_j(x), h)_H \\ &+ \sum_{j=1}^m \frac{\partial f}{\partial u_j} \Big|_{u_j=U_j(x)} (U''_j(x)h, h)_H. \end{aligned}$$

By (1),

$$\begin{aligned} \Delta_L F(x) &= \sum_{i,j=1}^m \frac{\partial^2 f}{\partial u_i \partial u_j} \Big|_{u_l=U_l(x)} \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n (U'_i(x), f_k)_H (U'_j(x), f_k)_H \\ &+ \sum_{j=1}^m \frac{\partial f}{\partial u_j} \Big|_{u_j=U_j(x)} \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n (U''_j(x) f_k, f_k)_H. \end{aligned}$$

But

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n (U'_i(x), f_k)_H (U'_j(x), f_k)_H = 0,$$

(since $(U'_i(x), f_k)_H \rightarrow 0$ as $k \rightarrow \infty$), and

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n (U''_j(x) f_k, f_k)_H = \Delta_L U_j(x).$$

From this we obtain

$$\Delta_L F(x) = \sum_{j=1}^m \frac{\partial f}{\partial u_j} \Big|_{u_j=U_j(x)} \Delta_L U_j(x).$$

Let Ω be a bounded domain in the Hilbert space H , that is, a bounded open set in H , and let $\bar{\Omega} = \Omega \cup \Gamma$ be the corresponding domain in H with boundary Γ .

In the space H we define a domain Ω with the surface Γ as follows:

$$\Omega = \{x \in H : 0 \leq Q(x) < R^2\}, \quad \Gamma = \{x \in H : Q(x) = R^2\},$$

where $Q(x)$ is a twice continuously differentiable function such that $\Delta_L Q(x) = \gamma$ for a nonzero positive constant γ . Domains of this type are called fundamental domains.

Let us give some examples of fundamental domains.

1) The ball $\bar{\Omega} = \{x \in H : \|x\|_H^2 \leq R^2\}$.

2) The ellipsoid $\bar{\Omega} = \{x \in H : (Bx, x)_H \leq R^2\}$, where $B = \gamma E + S(x)$, E is the identity operator and $S(x)$ is a compact operator on H .

Let us introduce a function $T(x) = \frac{R^2 - Q(x)}{\gamma}$. This function possesses the following properties:

$$\begin{aligned} 0 < T(x) &\leq \frac{R^2}{\gamma}, & \Delta_L T(x) &= -1 \quad \text{if } x \in \Omega, \\ T(x) &= 0 \quad \text{if } x \in \Gamma. \end{aligned}$$

3. THE PROBLEM WITHOUT INITIAL AND BOUNDARY CONDITIONS

Consider the nonlinear equation

$$(3) \quad f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0,$$

where $U(t, x)$ is a function on $[0, T] \times H$, $f(\xi, \eta, \zeta)$ is a given function with three arguments.

Theorem 1. 1. Let $f(\xi, \eta, \zeta)$ be a continuous twice differentiable function with three arguments taking values in the domain $\{U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\}$ in \mathbb{R}^3 .

2. Assume that one can solve the equation $f(\xi, \eta, c\eta) = 0$ with respect to η , $\eta = \phi(\xi, c)$ (although the original equation $f(\xi, \eta, \zeta) = 0$, in general, might have no solution η) and the solution admits separation of the variables ξ and c , i.e., $\phi(\xi, c) = \alpha(c)\beta(\xi)$ (for some functions $\alpha(c)$, $\beta(\xi)$ on \mathbb{R}^1 , $\beta(\xi) \neq 0$).

Then a solution of (3) can be given in an implicit form

$$(4) \quad \varphi(U(t, x)) = \alpha(\Psi(x))t + \delta(\Psi(x))\frac{\|x\|_H^2}{2} + \Phi(x),$$

where $\varphi(\xi) = \int \frac{d\xi}{\beta(\xi)}$, $\delta(c) = c\alpha(c)$, and $\Psi(x)$, $\Phi(x)$ are arbitrary harmonic functions on H .

Proof. We deduce from (4), using (2) and the relation $\delta(\Psi(x)) = \Psi(x)\alpha(\Psi(x))$, that

$$\begin{aligned} \varphi'_\xi(U(t, x))\frac{\partial U(t, x)}{\partial t} &= \alpha(\Psi(x)), \\ \varphi'_\xi(U(t, x))\Delta_L U(t, x) &= \alpha'_c(\Psi(x))\Delta_L \Psi(x)t + \Psi(x)\alpha'_c(\Psi(x))\Delta_L \Psi(x)\frac{\|x\|_H^2}{2} \\ &\quad + \Delta_L \Psi(x)\alpha(\Psi(x))\frac{\|x\|_H^2}{2} + \Psi(x)\alpha(\Psi(x))\frac{1}{2}\Delta_L \|x\|_H^2 + \Delta_L \Phi(x) \\ &= \Psi(x)\alpha(\Psi(x)) \end{aligned}$$

(since $\Delta_L \Psi(x) = \Delta_L \Phi(x) = 0$ by harmonicity and $\Delta_L \|x\|_H^2 = 2$ according to (1)).

Since $\varphi'_\xi = \frac{1}{\beta(\xi)}$ this implies that

$$\begin{aligned} \frac{\partial U(t, x)}{\partial t} &= \alpha(\Psi(x))\beta(U(t, x)) = \phi(U(t, x), \Psi(x)), \\ \Delta_L U(t, x) &= \delta(\Psi(x))\beta(U(t, x)) = \Psi(x)\phi(U(t, x), \Psi(x)). \end{aligned}$$

Substituting these relations into (3) we obtain

$$f\left(U(t, x), \phi(U(t, x), \Psi(x)), \Psi(x)\phi(U(t, x), \Psi(x))\right) = 0.$$

Due to condition 2) in the statement of the theorem this yields the identity

$$\phi(U(t, x), \Psi(x)) = \phi(U(t, x), \Psi(x)).$$

□

4. BOUNDARY PROBLEM

Consider the boundary problem

$$(5) \quad f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0 \quad \text{in } \Omega,$$

$$(6) \quad U(t, x) = G(t, x) \quad \text{on } \Gamma,$$

where $U(t, x)$ is a function on $[0, T] \times H$, $f(\xi, \eta, \zeta)$ is a given function having three arguments and $G(t, x)$ is a given function.

Theorem 2. 1. Let $f(\xi, \eta, \zeta)$ be a continuous twice differentiable function with three arguments defined in the domain $\{U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\}$ in \mathbb{R}^3 .

2. Assume that one can solve the equation $f(\xi, \eta, c\eta) = 0$ with respect to η , $\eta = \phi(\xi, c)$, and the solution admits the separation of variables ξ and c that is $\phi(\xi, c) = \alpha(c)\beta(\xi)$ (for some functions $\alpha(c)$, $\beta(\xi)$ on \mathbb{R}^1 , $\beta(\xi) \neq 0$).

3. Assume that there exist a primitive $\varphi(\xi) = \int \frac{d\xi}{\beta(\xi)}$ and its inverse function φ^{-1} .

4. Assume that the domain $\bar{\Omega}$ is fundamental.

5. Assume that in some functional space there exists a solution $V(\tau, x)$ of the boundary problem for the heat equation

$$(7) \quad \frac{\partial V(\tau, x)}{\partial \tau} = \Delta_L V(\tau, x) \quad \text{in } \Omega, \quad V(\tau, x)\Big|_{\Gamma} = G(\tau, x).$$

6. Consider the equation

$$(8) \quad \alpha'_c \left(\alpha^{-1} \left(\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=X+T(x)} \right) \right) [t - X] - \delta'_c \left(\alpha^{-1} \left(\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=X+T(x)} \right) \right) T(x) = 0,$$

where $\delta(c) = c\alpha(c)$, and assume that it can be solved with respect to $X = \chi(t, x)$, and $\chi(t, x)\Big|_{\Gamma} = t$.

Then the solution of the boundary problem (5), (6) in the same functional space is given by the formula

$$(9) \quad \varphi(U(t, x)) = \alpha(\psi(\chi(t, x)))[t - \chi(t, x)] - \delta(\psi(\chi(t, x)))T(x) + \varphi(V(\chi(t, x) + T(x), x)),$$

where

$$(10) \quad \psi(\chi(t, x)) = \alpha^{-1} \left(\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)} \right)$$

($\psi(z)$ is a function on \mathbb{R}^1).

Proof. Since $\varphi'_\xi = \frac{1}{\beta(\xi)}$ we deduce from (9)

$$\begin{aligned}
\varphi'_\xi(U(t, x)) \frac{\partial U(t, x)}{\partial t} &= \frac{1}{\beta(U(t, x))} \frac{\partial U(t, x)}{\partial t} \\
&= \alpha(\psi(\chi(t, x))) - \alpha(\psi(\chi(t, x))) \frac{\partial \chi(t, x)}{\partial t} \\
&\quad + \alpha'_c(\psi(\chi(t, x))) \psi'_z(\chi(t, x)) \frac{\partial \chi(t, x)}{\partial t} [t - \chi(t, x)] \\
&\quad - \delta'_c(\psi(\chi(t, x))) \psi'_z(\chi(t, x)) \frac{\partial \chi(t, x)}{\partial t} T(x) \\
&\quad + \frac{\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x) + T(x), x))} \frac{\partial \chi(t, x)}{\partial t} \\
&= \alpha(\psi(\chi(t, x))) + \left\{ \alpha'_c(\psi(\chi(t, x))) [t - \chi(t, x)] \right. \\
&\quad \left. - \delta'_c(\psi(\chi(t, x))) T(x) \right\} \psi'_z(\chi(t, x)) \frac{\partial \chi(t, x)}{\partial t} \\
&\quad - \left[\alpha(\psi(\chi(t, x))) - \frac{\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x) + T(x), x))} \right] \frac{\partial \chi(t, x)}{\partial t}.
\end{aligned}$$

Note that $\chi(t, x)$ solves (8) and due to (10) we have

$$(11) \quad \frac{\partial U(t, x)}{\partial t} = \alpha(\psi(\chi(t, x))) \beta(U(t, x)).$$

Since $\varphi'_\xi = \frac{1}{\beta(\xi)}$, and $\Delta_L T(x) = -1$, applying (2) we deduce from (9)

$$\begin{aligned}
\varphi'_\xi(U(t, x)) \Delta_L U(t, x) &= \frac{1}{\beta(U(t, x))} \Delta_L U(t, x) \\
&= -\alpha(\psi(\chi(t, x))) \Delta_L \chi(t, x) + \alpha'_c(\psi(\chi(t, x))) \psi'_z(\chi(t, x)) \Delta_L \chi(t, x) [t - \chi(t, x)] \\
&\quad - \delta'_c(\psi(\chi(t, x))) \psi'_z(\chi(t, x)) \Delta_L \chi(t, x) T(x) - \delta(\psi(\chi(t, x))) \Delta_L T(x) \\
&\quad + \frac{\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x) + T(x), x))} \left[\Delta_L \chi(t, x) + \Delta_L T(x) \right] + \frac{\Delta_L V(\tau, x) \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x) + T(x), x))} \\
&= \delta(\psi(\chi(t, x))) + \left\{ \alpha'_c(\psi(\chi(t, x))) [t - \chi(t, x)] \right. \\
&\quad \left. - \delta'_c(\psi(\chi(t, x))) T(x) \right\} \psi'_z(\chi(t, x)) \Delta_L \chi(t, x) \\
&\quad - \left[\alpha(\psi(\chi(t, x))) - \frac{\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x) + T(x), x))} \right] \Delta_L \chi(t, x) \\
&\quad - \frac{\left[\frac{\partial V(\tau, x)}{\partial \tau} - \Delta_L V(\tau, x) \right] \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x) + T(x), x))}.
\end{aligned}$$

Recall that $\chi(t, x)$ solves (8), $\frac{\partial V(\tau, x)}{\partial \tau} = \Delta_L V(\tau, x)$, then taking into account (10) we obtain

$$(12) \quad \Delta_L U(t, x) = \delta(\psi(\chi(t, x))) \beta(U(t, x)).$$

Substituting (11) and (12) into (5) we derive

$$(13) \quad f\left(U(t, x), \alpha(\psi(\chi(t, x))) \beta(U(t, x)), \psi(\chi(t, x)) \alpha(\psi(\chi(t, x))) \beta(U(t, x))\right) = 0.$$

By condition 2) in the statement of the theorem the identity

$$\alpha(\psi(\chi(t, x)))\beta(U(t, x)) = \alpha(\psi(\chi(t, x)))\beta(U(t, x))$$

can be deduced from (13).

At the surface Γ we have $T(x) = 0$, and $\chi(t, x) = t$. Setting $T(x) = 0$, $\chi(t, x) = t$ in (9) and keeping in mind that $V(t, x)|_{\Gamma} = G(t, x)$ we obtain

$$\varphi\left(U(t, x)\Big|_{\Gamma}\right) = \varphi\left(V(t, x)\Big|_{\Gamma}\right) = \varphi(G(t, x)) \quad \text{and} \quad U(t, x)\Big|_{\Gamma} = G(t, x).$$

□

5. INITIAL-BOUNDARY VALUE PROBLEM

Consider an initial-boundary value problem with uniform initial data

$$(14) \quad f\left(U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\right) = 0 \quad \text{in } \Omega,$$

$$(15) \quad U(0, x) = 0,$$

$$(16) \quad U(t, x) = G(t, x) \quad \text{on } \Gamma,$$

where $U(t, x)$ is a function on $[0, T] \times H$, $f(\xi, \eta, \zeta)$ is a given function with three arguments, $G(t, x)$ is a given function.

Theorem 3. 1. Let $f(\xi, \eta, \zeta)$ be a continuous twice differentiable function with three arguments defined in the domain $\{U(t, x), \frac{\partial U(t, x)}{\partial t}, \Delta_L U(t, x)\}$ in R^3 .

2. Assume that one can solve the equation $f(\xi, \eta, c\eta) = 0$ with respect to η , $\eta = \phi(\xi, c)$, and the solution admits the separation of variables ξ and c that is $\phi(\xi, c) = \alpha(c)\beta(\xi)$ (for some functions $\alpha(c), \beta(\xi)$ on $R^1, \beta(\xi) \neq 0$).

3. Assume that there exist a primitive $\varphi(\xi) = \int \frac{d\xi}{\beta(\xi)}$ and its inverse function φ^{-1} .

4. Assume that the domain $\bar{\Omega}$ is fundamental.

5. Assume that in some functional space there exists a solution $V(\tau, x)$ of the initial-boundary value problem for the heat equation

$$(17) \quad \frac{\partial V(\tau, x)}{\partial \tau} = \Delta_L V(\tau, x) \quad \text{in } \Omega, \quad V(0, x) = 0, \quad V(\tau, x)\Big|_{\Gamma} = G(t, x).$$

6. Assume that the equation

$$(18) \quad \alpha'_c \left(\alpha^{-1} \left(\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=X+T(x)} \right) \right) [t-X] - \delta'_c \left(\alpha^{-1} \left(\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=X+T(x)} \right) \right) T(x) = 0,$$

where $\delta(c) = \alpha(c)$, can be solved with respect to $X = \chi(t, x)$, and $\chi(t, x)\Big|_{\Gamma} = t$, $\chi(0, x) < r$.

7. Assume in addition that the function $G(t, x)$ is uniformly continuous in $\bar{\Omega}$ for each $t \in [0, T]$, and the mean value $\mathfrak{M}G(t, x + \sqrt{2T(x)}y)$ for all $t \in [0, T]$ and $G(t, x) = 0$, $G'_t(t, x) = 0$ for $t \leq r$ ($r > 0$) exists.

Then the solution of initial-boundary value problem (14)–(16) in this functional space is given by

$$(19) \quad \varphi(U(t, x)) = \alpha(\psi(\chi(t, x)))[t - \chi(t, x)] - \delta(\psi(\chi(t, x)))T(x) + \varphi(V(\chi(t, x) + T(x), x)),$$

where

$$(20) \quad \psi(\chi(t, x)) = \alpha^{-1} \left(\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)} \right) / \beta(V(\chi(t, x) + T(x), x))$$

($\psi(z)$ is a function on \mathbb{R}^1).

Proof. We can prove that the function given by (19) satisfies (14) in Ω and at the surface Γ $U(t, x) = G(t, x)$, in the similar to the proof of theorem 2.

Let us show that $U(0, x) = 0$.

First we prove that if $G(\tau, x) = 0$ for $\tau \leq 0$, then under the conditions of the theorem (namely condition (7)) the solution of the problem

$$\frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x) \quad \text{in } \Omega, \quad V(0, x) = 0, \quad V(t, x) \Big|_{\Gamma} = G(t, x)$$

can be written in the form

$$(21) \quad V(t, x) = \mathfrak{M}G(t - T(x), x + \sqrt{2T(x)}y),$$

where $\mathfrak{M}\Phi$ is a mean value of the function $\Phi(y)$ over the sphere $\|y\|_H^2 = 1$.

Eventually, at one hand

$$(22) \quad \frac{\partial V(t, x)}{\partial t} = \frac{\partial \mathfrak{M}G(t - T(x), x + \sqrt{2T(x)}y)}{\partial \tau}.$$

At the other hand using (2) we derive

$$\begin{aligned} \Delta_L V(t, x) &= - \frac{\partial \mathfrak{M}G(t - T(x), x + \sqrt{2T(x)}y)}{\partial \tau} \Delta_L T(x) + \Delta_L \mathfrak{M}G(\tau, x + \sqrt{2T(x)}y) \Big|_{\tau=t-T(x)}. \end{aligned}$$

It was shown in the paper [4] by E. M. Polischuk that if the function $F(x)$ is uniformly continuous in Ω and has a mean value $\mathfrak{M}F(x + \sqrt{2T(x)}y)$, then this mean value is a harmonic function in Ω , that is $\Delta_L \mathfrak{M}F(x + \sqrt{2T(x)}y) = 0$ ($x \in \Omega$).

Hence taking into account that $\Delta_L T(x) = -1$, we obtain

$$(23) \quad \Delta_L V(t, x) = \frac{\partial \mathfrak{M}G(t - T(x), x + \sqrt{2T(x)}y)}{\partial \tau}.$$

Substituting (22) and (23) into the equation $\frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x)$ we get the identity

$$\frac{\partial \mathfrak{M}G(t - T(x), x + \sqrt{2T(x)}y)}{\partial \tau} = \frac{\partial \mathfrak{M}G(t - T(x), x + \sqrt{2T(x)}y)}{\partial \tau}.$$

Setting $t = 0$ in (21) we obtain $V(0, x) = \mathfrak{M}G(-T(x), x + \sqrt{2T(x)}y) = 0$, since by conditions of the theorem we have $G(\tau, x) = 0$ for $\tau \leq 0$.

At the surface Γ $T(x) = 0$, and (21) yields $V(t, x) \Big|_{\Gamma} = \mathfrak{M}G(t, x) = G(t, x)$.

It results from (21) that

$$V(\chi(t, x) + T(x), x) = \mathfrak{M}G(\chi(t, x), x + \sqrt{2T(x)}y)$$

and thus

$$(24) \quad V(\chi(0, x) + T(x), x) = \mathfrak{M}G(\chi(0, x), x + \sqrt{2T(x)}y) = 0$$

(since by theorem conditions $\chi(0, x) \leq r$, and $G(\tau, x) = 0$ for $\tau \leq r$).

The relation

$$\alpha(\psi(\chi(t, x))) = \frac{\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)}}{\beta(V(\chi(t, x)) + T(x), x)} = \frac{\mathfrak{M}G'_\tau(\chi(t, x), x + \sqrt{2T(x)}y)}{\beta(V(\chi(t, x)) + T(x), x)}$$

follows from (20) and leads to $\alpha(\psi(\chi(0, x))) = 0$ (since by theorem conditions $\chi(0, x) < r$, and $G'(\tau, x) = 0$ for $\tau \leq r$).

Setting $t = 0$ in (19) and taking into account (24) and the existence of φ^{-1} , we deduce $\varphi(U(0, x)) = \varphi(V(\chi(0, t) + T(x), x)) = \varphi(0)$, that yields $U(0, x) = 0$. \square

Remark. If there exist unique solutions η , c and $\chi(t, x)$ to the equations $f(\xi, \eta, c\eta) = 0$, $\alpha(c) = z$ and (18) respectively, and if the initial-boundary value problem (17) for the heat equation has a unique solution in a certain functional class then the solution of the initial-boundary value problem (14)–(16) is unique in this functional space.

Example. Let us solve the initial-boundary value problem in a ball of the Hilbert space H : $\bar{\Omega} = \{x \in H : \|x\|_H^2 \leq R^2\}$

$$(25) \quad \left(\frac{\partial U(t, x)}{\partial t}\right)^3 - U(t, x) \left(\frac{\partial U(t, x)}{\partial t}\right)^2 + \left(\frac{\partial U(t, x)}{\partial t}\right) \Delta_L U(t, x) \\ = U(t, x) \Delta_L U(t, x) \quad \text{in } \Omega,$$

$$(26) \quad U(0, x) = 0,$$

$$(27) \quad U(t, x) \Big|_{\|x\|_H^2 = R^2} = g\left(t - \frac{1}{2}\|x\|_H^2\right),$$

where $g(\lambda) = \lambda^2$ if $\lambda \geq 0$, $g(\lambda) = 0$ if $\lambda \leq 0$.

To apply theorem 3 we note that the function f in (25) has the form

$$f(\xi, \eta, \zeta) = \eta^3 - \xi\eta^2 + \eta\zeta - \xi\zeta.$$

Hence the equation that appears in the condition 2) of theorem 3 has the form

$$(28) \quad \eta^3 - \xi\eta^2 + c\eta^2 - c\xi\eta = 0.$$

The solutions of the equation (28) have the form $\eta = -c$, $\eta = \xi$ and $\eta = 0$.

Let us take the solution $\eta = -c$. In this case $\eta = \phi(\xi, c) = -c$. Hence $\alpha(c) = -c$, $\beta(\xi) = 1$ which yields $\delta(c) = -c^2$, $\varphi(\xi) = \xi$.

A solution of the initial-boundary value problem for the heat equation

$$\frac{\partial V(\tau, x)}{\partial \tau} = \Delta_L V(\tau, x) \quad \text{in } \Omega, \quad V(0, x) = 0, \quad V(\tau, x) \Big|_{\|x\|_H^2 = R^2} = g\left(\tau - \frac{1}{2}\|x\|_H^2\right)$$

is given by the formula

$$V(\tau, x) = g\left(\tau + \frac{1}{2}\|x\|_H^2 - R^2\right).$$

It results

$$V(\tau, x) \Big|_{\tau=X+T} = g\left(X - \frac{R^2}{2}\right), \quad \frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=X+T} = 2\sqrt{g\left(X - \frac{R^2}{2}\right)}.$$

Now the equation (18) (in the condition 6) of theorem 3) takes the form

$$t - (1 - 4T(x))X - 2R^2T(x) = 0.$$

A solution of this equation is given by

$$X = \chi(t, x) = \frac{t - 2R^2T(x)}{1 - 4T(x)},$$

and in addition we have $\chi(t, x) \Big|_{\Gamma} = t$.

Since for such a form of $\chi(t, x)$

$$V(\chi(t, x) + T(x), x) = g\left(\chi(t, x) - \frac{R^2}{2}\right) = \frac{g\left(t - \frac{R^2}{2}\right)}{(1 - 4T(x))^2},$$

$$\frac{\partial V(\tau, x)}{\partial \tau} \Big|_{\tau=\chi(t, x)+T(x)} = \frac{2\sqrt{g\left(t - \frac{R^2}{2}\right)}}{1 - 4T(x)},$$

we deduce

$$\alpha(\psi(\chi(t, x))) = \frac{2\sqrt{g(t - R^2/2)}}{1 - 4T(x)}, \quad \delta(\psi(\chi(t, x))) = -\frac{4g(t - \frac{R^2}{2})}{(1 - 4T(x))^2}.$$

It results from (19) in theorem 3 that the solution of the problem (25)–(27) has the form

$$U(t, x) = \frac{g\left(t - \frac{R^2}{2}\right)}{1 - 2(R^2 - \|x\|_H^2)}.$$

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