

## QUASILINEAR PARABOLIC EQUATIONS WITH A LÉVY LAPLACIAN FOR FUNCTIONS OF INFINITE NUMBER OF VARIABLES

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*To dear Myroslav Gorbachuk on his 70th birthday.*

ABSTRACT. We construct solutions to initial, boundary and initial-boundary value problems for quasilinear parabolic equations with an infinite dimensional Lévy Laplacian  $\Delta_L$ ,

$$\frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) + f_0(U(t, x)),$$

in fundamental domains of a Hilbert space. The solution is defined in the functional class where a solution of the corresponding problem for the heat equation  $\frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x)$  exists.

### 1. INTRODUCTION

In 1919 P. Lévy considered a quasilinear elliptic equation

$$\Delta_L U(x) = f(U(x)),$$

where  $U(x)$  is a function defined on a Hilbert space  $H$  and  $f(\xi)$  is a function of a scalar argument. P. Lévy showed in [1] that a general solution of this equation is given implicitly by the relation

$$\varphi(U(x)) - \frac{1}{2} \|x\|_H^2 = \Psi(x),$$

where  $\varphi(\xi) = \int \frac{d\xi}{f(\xi)}$  and  $\Psi(x)$  is an arbitrary harmonic function, and reduced solution of the Dirichlet problem for this equation in a bounded domain of the Hilbert space to the Dirichlet problem for the Lévy Laplace equation. Later these problems were described in the book by P. Lévy [2].

Solution of the Cauchy problem for a quasilinear parabolic equation with the Lévy Laplacian were constructed by M. Feller in paper [3]. One can find more references in the book by M. Feller [4].

In this article we construct a solution of a boundary value problem and an initial-boundary value problem for quasilinear parabolic equations with the Lévy Laplacian

$$\frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) + f_0(U(t, x)),$$

(here  $f_0(\xi)$  is a function defined on  $\mathbb{R}^1$ ) for a fundamental domain in a real infinite dimensional Hilbert space. To make the article self-consistent we give the expression of the solution of the Cauchy problem for a quasilinear parabolic equation with the Lévy Laplacian.

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## 2. PRELIMINARIES

Let  $H$  be a real infinite dimensional Hilbert space.

The infinite dimensional Laplacian defined by P. Lévy can be described as follows. If a scalar function  $F$  defined on  $H$  is twice strongly differentiable at a point  $x_0$  then the Lévy Laplacian of  $F$  at the point  $x_0$  is defined (if it exists) by the formula

$$(1) \quad \Delta_L F(x_0) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n (F''(x_0) f_k, f_k)_H,$$

where  $F''(x)$  is the Hessian of  $F(x)$ , and  $\{f_k\}_1^\infty$  is an orthonormal basis in  $H$ .

We shall recall an important property of the Lévy Laplacian described in [2] that we will need in the sequel. Assume that

$$F(x) = f(U_1(x), \dots, U_m(x)),$$

where  $f(u_1, \dots, u_m)$  is a twice continuously differentiable function of  $m$  variables in a domain  $\{U_1(x), \dots, U_m(x)\} \subset \mathbb{R}^m$ , where  $(U_1(x), \dots, U_m(x))$  is a vector of values of functions  $U_1(x), \dots, U_m(x)$ . We assume that  $U_j(x)$  are uniformly continuous in a bounded domain  $\Omega \subset H$  and strongly twice differentiable functions, and  $\Delta_L U_j(x)$  exist ( $j = 1, \dots, m$ ). Then  $\Delta_L F(x)$  exists and

$$(2) \quad \Delta_L F(x) = \sum_{j=1}^m \frac{\partial f(u)}{\partial u_j} \Big|_{u_j=U_j(x)} \Delta_L U_j(x).$$

We deduce two consequences from (2).

1) If the functions  $U_k(x)$  are harmonic in some domain  $\Omega$  ( $k = 1, \dots, m$ ), then the function  $F(x)$  is also harmonic in  $\Omega$ .

2) The Lévy Laplacian is a "derivative". Namely, for  $F(x) = U_1(x)U_2(x)$  we have  $\Delta_L[U_1(x)U_2(x)] = \Delta_L U_1(x) \cdot U_2(x) + U_1(x) \cdot \Delta_L U_2(x)$ .

Let  $\Omega$  be a bounded domain in the Hilbert space  $H$  (that is a bounded open set in  $H$ ), and  $\bar{\Omega} = \Omega \cup \Gamma$  be a domain in  $H$  with boundary  $\Gamma$ .

Consider the domain  $\Omega$  in  $H$  with boundary  $\Gamma$  of the form

$$\Omega = \{x \in H : 0 \leq Q(x) < R^2\}, \quad \Gamma = \{x \in H : Q(x) = R^2\},$$

where  $Q(x)$  is a twice strongly differentiable function such that  $\Delta_L Q(x) = \gamma$ , where  $\gamma > 0$  is a positive constant. Such domains are called fundamental domains.

Examples of fundamental domains are

- 1) a ball  $\bar{\Omega} = \{x \in H : \|x\|_H^2 \leq R^2\}$ ,
- 2) an ellipsoid  $\bar{\Omega} = \{x \in H : (Bx, x)_H \leq R^2\}$ , where  $B = \gamma E + S(x)$ ,  $E$  is the identity operator and  $S(x)$  is a compact operator acting in  $H$ .

Consider the function

$$T(x) = \frac{R^2 - Q(x)}{\gamma},$$

that will be used in the sequel. The function  $T(x)$  possesses the following properties:

$$0 < T(x) \leq \frac{R^2}{\gamma}, \quad \Delta_L T(x) = -1 \quad \text{if } x \in \Omega,$$

$$T(x) = 0 \quad \text{if } x \in \Gamma.$$

3. THE CAUCHY PROBLEM

Consider the Cauchy problem

$$(3) \quad \frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) + f_0(U(t, x)),$$

$$(4) \quad U(0, x) = U_0(x),$$

where  $U(t, x)$  is a function on  $[0, \mathfrak{T}] \times H$ ,  $f_0(\xi)$  is a given function of one variable,  $U_0(x)$  is a given function defined on  $H$ .

**Theorem 1.** *Let  $f_0(\xi)$  be a differentiable function in the domain  $\{U(t, x)\}$  ( $\{U(t, x)\}$  is a domain in  $\mathbb{R}^1$ , where the function  $U(t, x)$  takes its values).*

*Assume that there exists a primitive  $\varphi(\xi) = \int \frac{d\xi}{f_0(\xi)}$  and the inverse function  $\varphi^{-1}$ .*

*Assume in addition that in a certain functional class  $\mathcal{F}$  there exists a solution of the Cauchy problem for the heat equation*

$$(5) \quad \frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x), \quad V(0, x) = U_0(x).$$

*Then the solution  $U(t, x)$  of the Cauchy problem (3), (4) in the same class  $\mathcal{F}$  is given by the equation*

$$(6) \quad \varphi(U(t, x)) = t + \varphi(V(t, x)),$$

*i.e.,  $U(t, x) = \varphi^{-1}(t + \varphi(V(t, x)))$ .*

*Proof.* We deduce from (6) that

$$\varphi'_\xi(U(t, x)) \frac{\partial U(t, x)}{\partial t} = 1 + \varphi'_\xi(V(t, x)) \frac{\partial V(t, x)}{\partial t}.$$

For  $m = 1$ , we deduce from (6) using (2) that

$$\varphi'_\xi(U(t, x)) \Delta_L U(t, x) = \varphi'_\xi(V(t, x)) \Delta_L V(t, x).$$

But  $\varphi'_\xi(\xi) = \frac{1}{f_0(\xi)}$ , hence

$$(7) \quad \frac{\partial U(t, x)}{\partial t} = f_0(U(t, x)) + \frac{f_0(U(t, x))}{f_0(V(t, x))} \frac{\partial V(t, x)}{\partial t},$$

$$(8) \quad \Delta_L U(t, x) = \frac{f_0(U(t, x))}{f_0(V(t, x))} \Delta_L V(t, x).$$

Substituting (7) and (8) into (3), we derive

$$f_0(U(t, x)) + \frac{f_0(U(t, x))}{f_0(V(t, x))} \frac{\partial V(t, x)}{\partial t} = \frac{f_0(U(t, x))}{f_0(V(t, x))} \Delta_L V(t, x) + f_0(U(t, x)),$$

i.e., (6) satisfies equation (3).

Choosing  $t = 0$  in (6) and taking into account that  $V(0, x) = U_0(x)$ , we deduce that  $U(0, x) = U_0(x)$ . □

*Example 1.* Let us construct a solution of the Cauchy problem

$$(9) \quad \frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) - U^3(t, x),$$

$$(10) \quad U(0, x) = h\left(\frac{1}{2}\|x\|_H^2\right),$$

where  $h(\lambda)$  is a smooth function on  $(-\infty < \lambda < \infty)$ .

Note that in (9)  $f_0(\xi) = -\xi^3$  that yields  $\varphi(\xi) = -\int \frac{d\xi}{\xi^3} = \frac{1}{2\xi^2}$ , and hence  $\varphi^{-1}(z) = \frac{1}{\sqrt{2z}}$ .

A solution of the Cauchy problem for the heat equation

$$\frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x), \quad V(0, x) = h\left(\frac{1}{2}\|x\|_H^2\right)$$

has the form

$$V(t, x) = h\left(t + \frac{1}{2}\|x\|_H^2\right).$$

Thus, by (6) we obtain a solution of the problem (9), (10) in the form

$$U(t, x) = \frac{1}{\sqrt{2\left[t + \frac{1}{2}[h(t + \frac{1}{2}\|x\|_H^2)]^{-2}\right]}}.$$

#### 4. BOUNDARY VALUE PROBLEM

Consider a boundary value problem

$$(11) \quad \frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) + f_0(U(t, x)) \quad \text{in } \Omega,$$

$$(12) \quad U(t, x) = G(t, x) \quad \text{on } \Gamma,$$

where  $U(t, x)$  is a function on  $[0, \mathfrak{T}] \times H$ ,  $f_0(\xi)$  is a given function of one dimensional argument and  $G(t, x)$  is a given function.

**Theorem 2.** *Let  $f_0(\xi)$  be a differentiable function in the domain  $\{U(t, x)\}$ .*

*Let there exist both the primitive  $\varphi(\xi) = \int \frac{d\xi}{f_0(\xi)}$  and its inverse function  $\varphi^{-1}$ .*

*Let, in addition,  $\bar{\Omega}$  be a fundamental domain.*

*Assume that in a certain functional class  $\mathcal{F}$  there exists a solution of the boundary value problem for the heat equation,*

$$(13) \quad \frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x) \quad \text{in } \Omega, \quad V(t, x)|_{\Gamma} = G(t, x).$$

*Then, in this class  $\mathcal{F}$ , a solution of the boundary value problem (11), (12) is defined by the formula*

$$(14) \quad \varphi(U(t, x)) = T(x) + \varphi(V(t, x)),$$

*i.e.,  $U(t, x) = \varphi^{-1}(T(x) + \varphi(V(t, x)))$ . Here  $T(x) = \frac{R^2 - Q(x)}{\gamma}$  (see p. 2).*

*Proof.* We deduce from (14) the relation

$$\varphi'_\xi(U(t, x)) \frac{\partial U(t, x)}{\partial t} = \varphi'_\xi(V(t, x)) \frac{\partial V(t, x)}{\partial t}.$$

For  $m = 1$ , we deduce from (14) using (2) the relation

$$\varphi'_\xi(U(t, x)) \Delta_L U(t, x) = \Delta_L T(x) + \varphi'_\xi(V(t, x)) \Delta_L V(t, x).$$

Since  $\Delta_L T(x) = -1$  and  $\varphi'_\xi(\xi) = \frac{1}{f_0(\xi)}$ , we get

$$(15) \quad \frac{\partial U(t, x)}{\partial t} = \frac{f_0(U(t, x))}{f_0(V(t, x))} \frac{\partial V(t, x)}{\partial t},$$

$$(16) \quad \Delta_L U(t, x) = -f_0(U(t, x)) + \frac{f_0(U(t, x))}{f_0(V(t, x))} \Delta_L V(t, x).$$

Substituting (15) and (16) into (11), we derive

$$\frac{f_0(U(t, x))}{f_0(V(t, x))} \frac{\partial V(t, x)}{\partial t} = -f_0(U(t, x)) + \frac{f_0(U(t, x))}{f_0(V(t, x))} \Delta_L V(t, x) + f_0(U(t, x)),$$

i.e., (14) satisfies equation (11).

At the surface  $\Gamma$  we obtain  $T(x) = 0$ . Substituting  $T(x) = 0$  into (14) and keeping in mind that  $V(t, x)|_{\Gamma} = G(t, x)$ , we deduce the equality  $U(t, x)|_{\Gamma} = G(t, x)$ .  $\square$

*Example 2.* Let us construct a solution of the boundary value problem without initial data in the ball of the space  $H$ ,  $\bar{\Omega} = \{x \in H : \|x\|_H^2 \leq R^2\}$ ,

$$(17) \quad \frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) - U^3(t, x) \quad \text{in } \Omega,$$

$$(18) \quad U(t, x)|_{\|x\|_H^2=R^2} = g\left(t - \frac{1}{2}\|x\|_H^2\right),$$

where  $g(\lambda)$  is a smooth function on  $(-\infty < \lambda < \infty)$ .

For the considered domain  $\Omega = \{x \in H : \|x\|_H^2 \leq R^2\}$ , we have  $T(x) = \frac{R^2 - \|x\|_H^2}{2}$ .

In the equation (17) we have  $f_0(\xi) = -\xi^3$ , which yields  $\varphi(\xi) = \frac{1}{2\xi^2}$ ,  $\varphi^{-1}(z) = \frac{1}{\sqrt{2z}}$ .

A solution of the boundary value problem without initial data for the heat equation

$$\frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x), \quad V(t, x)|_{\|x\|_H^2=R^2} = g\left(t - \frac{1}{2}\|x\|_H^2\right),$$

has the form

$$V(t, x) = g\left(t + \frac{1}{2}\|x\|_H^2 - R^2\right).$$

Hence by (14) we obtain a solution of the problem (17), (18) in the form

$$U(t, x) = \frac{1}{\sqrt{\left[R^2 - \|x\|_H^2 + \left[g\left(t + \frac{1}{2}\|x\|_H^2 - R^2\right)\right]^2\right]}}.$$

## 5. INITIAL-BOUNDARY VALUE PROBLEM

Consider the initial-boundary value problem

$$(19) \quad \frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) + f_0(U(t, x)) \quad \text{in } \Omega,$$

$$(20) \quad U(0, x) = U_0(x),$$

$$(21) \quad U(t, x) = G(t, x) \quad \text{on } \Gamma,$$

where  $U(t, x)$  is a function defined on  $[0, T] \times H$ ,  $f_0(\xi)$  is a given function of one variable and  $U_0(x), G(t, x)$  are given functions.

**Theorem 3.** *Let  $f_0(\xi)$  be a differentiable function in the domain  $\{U(t, x)\}$ .*

*Let there exist the primitive  $\varphi(\xi) = \int \frac{d\xi}{f_0(\xi)}$  and the inverse function  $\varphi^{-1}$ .*

*Let the domain  $\bar{\Omega}$  be fundamental.*

*Assume that, in a certain functional class  $\mathcal{F}$ , there exist solutions of the initial-boundary value problems*

$$(22) \quad \frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x) \quad \text{in } \Omega, \quad V(0, x) = U_0(x), \quad V(t, x)|_{\Gamma} = G(t, x).$$

*Then a solution of the initial-boundary value problem (19)–(21) in this class  $\mathcal{F}$  is given by the formula*

$$(23) \quad \varphi(U(t, x)) = \tau(t, T(x)) + \varphi(V(t, x)),$$

*i.e.,  $U(t, x) = \varphi^{-1}(\tau(t, T(x)) + \varphi(V(t, x)))$ , where  $\tau(t, T(x)) = t - q(t - T(x))$ ,  $q(\lambda) = \lambda$  if  $\lambda \geq 0$  and  $q(\lambda) = 0$  if  $\lambda \leq 0$ ,  $T(x) = \frac{R^2 - Q(x)}{\gamma}$ .*

If the initial-boundary value problem (22) possesses a unique solution in the same functional class then the solution of the initial-boundary value problem (19)–(21) is unique in this class.

*Proof.* We deduce from (23) that

$$\varphi'_\xi(U(t, x)) \frac{\partial U(t, x)}{\partial t} = \frac{\partial \tau(t, x)}{\partial t} + \varphi'_\xi(V(t, x)) \frac{\partial V(t, x)}{\partial t}.$$

For  $m = 1$ , we deduce from (23) using (2) that

$$\varphi'_\xi(U(t, x)) \Delta_L U(t, x) = \Delta_L \tau(t, x) + \varphi'_\xi(V(t, x)) \Delta_L V(t, x).$$

Recall that  $\varphi'_\xi(\xi) = \frac{1}{f_0(\xi)}$ , which yields

$$(24) \quad \frac{\partial U(t, x)}{\partial t} = f_0(U(t, x)) \frac{\partial \tau(t, x)}{\partial t} + \frac{f_0(U(t, x))}{f_0(V(t, x))} \frac{\partial V(t, x)}{\partial t},$$

$$(25) \quad \Delta_L U(t, x) = f_0(U(t, x)) \Delta_L \tau(t, x) + \frac{f_0(U(t, x))}{f_0(V(t, x))} \Delta_L V(t, x).$$

Substituting (24) and (25) into (19), we obtain

$$\begin{aligned} f_0(U(t, x)) \frac{\partial \tau(t, x)}{\partial t} + \frac{f_0(U(t, x))}{f_0(V(t, x))} \frac{\partial V(t, x)}{\partial t} \\ = f_0(U(t, x)) \Delta_L \tau(t, x) + \frac{f_0(U(t, x))}{f_0(V(t, x))} \Delta_L V(t, x) + f_0(U(t, x)), \end{aligned}$$

i.e., (23) satisfies equation (19) (since  $\frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x)$ ,  $\frac{\partial \tau(t, T)}{\partial t} = \Delta_L \tau(t, T) + 1$ ).

Setting  $t = 0$  in (23) and taking into account that  $V(0, x) = U_0(x)$ , and  $\tau(0, T) = 0$ , we obtain the equality  $U(0, x) = U_0(x)$ .

Since  $V(t, x) = G(t, x)$ ,  $\tau(t, T) = 0$  on the surface  $\Gamma$ , we deduce from (23) that  $U(t, x)|_\Gamma = G(t, x)$ .

The final statement of the theorem is obvious.  $\square$

*Example 3.* Let us construct a solution of the initial-boundary value problem in the ball  $\bar{\Omega} = \{x \in H : \|x\|_H^2 \leq R^2\}$  of the space  $H$ ,

$$(26) \quad \frac{\partial U(t, x)}{\partial t} = \Delta_L U(t, x) - U^3(t, x) \quad \text{in } \Omega,$$

$$(27) \quad U(0, x) = h\left(\frac{1}{2}\|x\|_H^2\right),$$

$$(28) \quad U(t, x)|_{\|x\|_H^2=R^2} = g\left(t - \frac{1}{2}\|x\|_H^2\right),$$

where  $h(\lambda)$  is a smooth function on the positive half axis with the support  $[0, \frac{R^2}{2}]$ , and  $g(\lambda)$  is a smooth function such that  $g(\lambda) = 0$  for  $\lambda \leq 0$  (note that comparing with Examples 1 and 2 we need some additional assumptions concerning the functions  $h(\lambda)$  and  $g(\lambda)$ ).

In the considered domain  $\bar{\Omega} = \{x \in H : \|x\|_H^2 \leq R^2\}$ , we have  $T(x) = \frac{R^2 - \|x\|_H^2}{2}$ .

In equation (26),  $f_0(\xi) = -\xi^3$  and, hence,  $\varphi(\xi) = -\int \frac{d\xi}{\xi^3} = \frac{1}{\xi^2}$  and  $\varphi^{-1}(z) = \frac{1}{\sqrt{2z}}$ .

Consider the initial-boundary value problem for the heat equation in the Shilov class of functions

$$\frac{\partial V(t, x)}{\partial t} = \Delta_L V(t, x) \quad \text{in } \Omega, \quad V(0, x) = h\left(\frac{1}{2}\|x\|_H^2\right),$$

$$(29) \quad V(t, x) \Big|_{\|x\|_H^2=R^2} = g\left(t - \frac{1}{2}\|x\|_H^2\right).$$

Its solution has the form

$$V(t, x) = h\left(t + \frac{1}{2}\|x\|_H^2\right) + g\left(t + \frac{1}{2}\|x\|_H^2 - R^2\right).$$

It results from (23) that the solution of the problem (26)–(28) has the form

$$U(t, x) = \frac{1}{\sqrt{2\left\{t - q\left(t + \frac{\|x\|_H^2}{2} - \frac{R^2}{2}\right) + \frac{1}{2}\left[h\left(t + \frac{1}{2}\|x\|_H^2\right) + g\left(t + \frac{1}{2}\|x\|_H^2 - R^2\right)\right]^{-2}\right\}}}.$$

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