

# SCHRÖDINGER OPERATORS WITH PURELY DISCRETE SPECTRUM

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*Dedicated to A. Ya. Povzner*

ABSTRACT. We prove that  $-\Delta + V$  has purely discrete spectrum if  $V \geq 0$  and, for all  $M$ ,  $|\{x \mid V(x) < M\}| < \infty$  and various extensions.

## 1. INTRODUCTION

Our main goal in this note is to explore one aspect of the study of Schrödinger operators

$$(1.1) \quad H = -\Delta + V$$

which we will suppose have  $V$ 's which are nonnegative and in  $L^1_{\text{loc}}(\mathbb{R}^\nu)$ , in which case (see, e.g., Simon [15])  $H$  can be defined as a form sum. We are interested here in criteria under which  $H$  has purely discrete spectrum, that is,  $\sigma_{\text{ess}}(H)$  is empty. This is well known to be equivalent to proving  $(H+1)^{-1}$  or  $e^{-sH}$  for any (and so all)  $s > 0$  is compact (see [9, Thm. XIII.16]). One of the most celebrated elementary results on Schrödinger operators is that this is true if

$$(1.2) \quad \lim_{|x| \rightarrow \infty} V(x) = \infty.$$

But (1.2) is not necessary. Simple examples where (1.2) fails but  $H$  still has compact resolvent were noted first by Rellich [10]—one of the most celebrated examples is in  $\nu = 2$ ,  $x = (x_1, x_2)$ , and

$$(1.3) \quad V(x_1, x_2) = x_1^2 x_2^2$$

where (1.2) fails in a neighborhood of the axes. For proof of this and discussions of eigenvalue asymptotics, see [11, 16, 17, 20, 21].

There are known necessary and sufficient conditions on  $V$  for discrete spectrum in terms of capacities of certain sets (see, e.g., Maz'ya [6]), but the criteria are not always so easy to check. Thus, I was struck by the following simple and elegant theorem:

**Theorem 1.** *Define*

$$(1.4) \quad \Omega_M(V) = \{x \mid 0 \leq V(x) < M\}.$$

*If (with  $|\cdot|$  Lebesgue measure)*

$$(1.5) \quad |\Omega_M(V)| < \infty$$

*for all  $M$ , then  $H$  has purely discrete spectrum.*

I learned of this result from Wang–Wu [25], but there is much related work. I found an elementary proof of Theorem 1 and decided to write it up as a suitable tribute and

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appreciation of A. Ya. Povzner, whose work on continuum eigenfunction expansions for Schrödinger operators in scattering situation [7] was seminal and inspired me as a graduate student forty years ago!

The proof has a natural abstraction:

**Theorem 2.** *Let  $\mu$  be a measure on a locally compact space,  $X$  with  $L^2(X, d\mu)$  separable. Let  $L_0$  be a selfadjoint operator on  $L^2(X, d\mu)$  so that its semigroup is ultracontractive ([1]): For some  $s > 0$ ,  $e^{-sL_0}$  maps  $L^2$  to  $L^\infty(X, d\mu)$ . Suppose  $V$  is a nonnegative multiplication operator so that*

$$(1.6) \quad \mu(\{x \mid 0 \leq V(x) < M\}) < \infty$$

for all  $M$ . Then  $L = L_0 + V$  has purely discrete spectrum.

*Remark.* By  $L_0 + V$ , we mean the operator obtained by applying the monotone convergence theorem for forms (see, e.g., [13, 14]) to  $L_0 + \min(V(x), k)$  as  $k \rightarrow \infty$ .

The reader may have noticed that (1.3) does not obey Theorem 1 (but, e.g.,

$$V(x_1, x_2) = x_1^2 x_2^4 + x_1^4 x_2^2$$

does). But our proof can be modified to a result that does include (1.3). Given a set  $\Omega$  in  $\mathbb{R}^\nu$ , define for any  $x$  and any  $\ell > 0$ ,

$$(1.7) \quad \omega_x^\ell(\Omega) = |\Omega \cap \{y \mid |y - x| \leq \ell\}|.$$

For example, for (1.3), for  $x \in \Omega_M$ ,

$$(1.8) \quad \omega_x^\ell(\Omega_M) \leq \frac{C_\ell}{|x| + 1}.$$

We will say a set  $\Omega$  is  $r$ -polynomially thin if

$$(1.9) \quad \int_{x \in \Omega} \omega_x^\ell(\Omega)^r d^\nu x < \infty$$

for all  $\ell$ . For the example in (1.3),  $\Omega_M$  is  $r$ -polynomially thin for any  $M$  and any  $r > 0$ . We'll prove

**Theorem 3.** *Let  $V$  be a nonnegative potential so that for any  $M$ , there is an  $r > 0$  so that  $\Omega_M$  is  $r$ -polynomially thin. Then  $H$  has purely discrete spectrum.*

As mentioned, this covers the example in (1.3). It is not hard to see that if  $P(x)$  is any polynomial in  $x_1, \dots, x_\nu$  so that for no  $v \in \mathbb{R}^\nu$  is  $\vec{v} \cdot \vec{\nabla} P \equiv 0$  (i.e.,  $P$  isn't a function of fewer than  $\nu$  linear variables), then  $V(x) = P(x)^2$  obeys the hypotheses of Theorem 3.

In Section 2, we'll present a simple compactness criterion on which all theorems rely. In Section 3, we'll prove Theorems 1 and 2. In Section 4, we'll prove Theorem 3.

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## 2. SEGAL'S LEMMA

Segal [12] proved the following result, sometimes called Segal's lemma:

**Proposition 2.1.** *For  $A, B$  positive selfadjoint operators,*

$$(2.1) \quad \|e^{-(A+B)}\| \leq \|e^{-A}e^{-B}\|.$$

*Remarks.* 1.  $A + B$  can always be defined as a closed quadratic form on  $Q(A) \cap Q(B)$ . That defines  $e^{-(A+B)}$  on  $\overline{Q(A) \cap Q(B)}$  and we set it to 0 on the orthogonal complement. Since the Trotter product formula is known in this generality (see Kato [5]), (2.1) holds in that generality.

2. Since  $\|C^*C\| = \|C\|^2$ ,  $\|e^{-A/2}e^{-B/2}\|^2 = \|e^{-B/2}e^{-A}e^{-B/2}\|$ , and since  
 $\|e^{-(A+B)/2}\|^2 = \|e^{-(A+B)}\|$ ,

(2.1) is equivalent to

$$(2.2) \quad \|e^{-A+B}\| \leq \|e^{-B/2}e^{-A}e^{-B/2}\|$$

which is the way Segal [12] stated it.

3. Somewhat earlier, Golden [4] and Thompson [22] proved

$$(2.3) \quad \mathrm{Tr}(e^{-(A+B)}) \leq \mathrm{Tr}(e^{-A}e^{-B})$$

and Thompson [23] later extended this to any symmetrically normed operator ideal.

*Proof.* There are many; see, for example, Simon [18, 19]. Here is the simplest, due to Deift [2, 3]: If  $\sigma$  is the spectrum of an operator

$$(2.4) \quad \sigma(CD) \setminus \{0\} = \sigma(DC) \setminus \{0\}$$

so with  $\sigma_r$  the spectral radius,

$$(2.5) \quad \sigma_r(CD) = \sigma_r(DC) \leq \|DC\|.$$

If  $CD$  is selfadjoint,  $\sigma_r(CD) = \|CD\|$ , so

$$(2.6) \quad CD \text{ selfadjoint} \Rightarrow \|CD\| \leq \|DC\|.$$

Thus,

$$(2.7) \quad \|e^{-A/2}e^{-B/2}\|^2 = \|e^{-B/2}e^{-A}e^{-B/2}\| \leq \|e^{-A}e^{-B}\|.$$

By induction,

$$(2.8) \quad \|(e^{-A/2^n}e^{-B/2^n})^{2^n}\| \leq \|e^{-A/2^n}e^{-B/2^n}\|^{2^n} \leq \|e^{-A}e^{-B}\|.$$

Take  $n \rightarrow \infty$  and use the Trotter product formula to get (2.1).  $\square$

In [18], I noted that this implies for any symmetrically normed trace ideal,  $\mathcal{I}_\Phi$ , that

$$(2.9) \quad e^{-A/2}e^{-B}e^{-A/2} \in \mathcal{I}_\Phi \Rightarrow e^{-(A+B)} \in \mathcal{I}_\Phi.$$

I explicitly excluded the case  $\mathcal{I}_\Phi = \mathcal{I}_\infty$  (the compact operators) because the argument there doesn't show that, but it is true—and the key to this paper!

Since

$$C \in \mathcal{I}_\infty \Leftrightarrow C^*C \in \mathcal{I}_\infty$$

and  $e^{-(A+B)} \in \mathcal{I}_\infty$  if and only if  $e^{-\frac{1}{2}(A+B)} \in \mathcal{I}_\infty$ , it doesn't matter if we use the symmetric form (2.2) or the following asymmetric form which is more convenient in applications.

**Theorem 2.2.** *Let  $\mathcal{I}_\infty$  be the ideal of compact operators on some Hilbert space,  $\mathcal{H}$ . Let  $A, B$  be nonnegative selfadjoint operators. Then*

$$(2.10) \quad e^{-A}e^{-B} \in \mathcal{I}_\infty \Rightarrow e^{-(A+B)} \in \mathcal{I}_\infty.$$

*Proof.* For any bounded operator,  $C$ , define  $\mu_n(C)$  by

$$(2.11) \quad \mu_n(C) = \min_{\psi_1 \dots \psi_{n-1}} \sup_{\substack{\|\varphi\|=1 \\ \varphi \perp \psi_1, \dots, \psi_{n-1}}} \|C\varphi\|.$$

By the min-max principle (see [9, Sect. XIII.1]),

$$(2.12) \quad \lim_{n \rightarrow \infty} \mu_n(C) = \sup(\sigma_{\mathrm{ess}}(|C|))$$

and  $\mu_n(C)$  are the singular values if  $C \in \mathcal{I}_\infty$ . In particular,

$$(2.13) \quad C \in \mathcal{I}_\infty \Leftrightarrow \lim_{n \rightarrow \infty} \mu_n(C) = 0.$$

Let  $\wedge^\ell(\mathcal{H})$  be the antisymmetric tensor product (see [8, Sects. II.4, VIII.10], [9, Sect. XIII.17], and [18, Sect. 1.5]). As usual (see [18, eqn. (1.14)]),

$$(2.14) \quad \|\wedge^m(C)\| = \prod_{j=1}^m \mu_j(C).$$

Since  $\mu_1 \geq \mu_2 \geq \dots \geq 0$ , we have

$$(2.15) \quad \lim_{n \rightarrow \infty} \mu_n(C) = \lim_{n \rightarrow \infty} (\mu_1(C) \dots \mu_n(C))^{1/n}.$$

(2.13)–(2.15) imply

$$(2.16) \quad C \in \mathcal{I}_\infty \Leftrightarrow \lim_{n \rightarrow \infty} \|\wedge^n(C)\|^{1/n} = 0.$$

As usual, there is a selfadjoint operator,  $d \wedge^n(A)$  on  $\wedge^n(\mathcal{H})$  so

$$(2.17) \quad \wedge^n(e^{-tA}) = e^{-t d \wedge^n(A)}$$

so Segal's lemma implies that

$$(2.18) \quad \|\wedge^n(e^{-(A+B)})\| \leq \|\wedge^n(e^{-A}) \wedge^n(e^{-B})\| = \|\wedge^n(e^{-A}e^{-B})\|.$$

Thus,

$$(2.19) \quad \lim_{n \rightarrow \infty} \|\wedge^n(e^{-(A+B)})\|^{1/n} \leq \lim_{n \rightarrow \infty} \|\wedge^n(e^{-A}e^{-B})\|^{1/n}.$$

By (2.16), we obtain (2.10).  $\square$

### 3. PROOFS OF THEOREMS 1 AND 2

*Proof of Theorem 1.* By Theorem 2.2, we need only show  $C = e^\Delta e^{-V}$  is compact. Write

$$(3.1) \quad C = C_m + D_m$$

where

$$(3.2) \quad C_m = C \chi_{\Omega_m}, \quad D_m = C \chi_{\Omega_m^c}$$

with  $\chi_S$  the operator of multiplication by the characteristic function of a set  $S \subset \mathbb{R}^\nu$ .

$$\|e^{-V} \chi_{\Omega_m^c}\|_\infty \leq e^{-m}$$

and  $\|e^\Delta\| = 1$ , so

$$(3.3) \quad \|D_m\| \leq e^{-m}$$

and thus,

$$(3.4) \quad \lim_{m \rightarrow \infty} \|C - C_m\| = 0.$$

If we show each  $C_m$  is compact, we are done. We know  $e^\Delta$  has integral kernel  $f(x-y)$  with  $f$  a Gaussian, so in  $L^2$ . Clearly, since  $V$  is positive,  $C_m$  has an integral kernel  $C_m(x, y)$  dominated by

$$(3.5) \quad |C_m(x, y)| \leq f(x-y) \chi_{\Omega_m}(y).$$

Thus,

$$\int |C_m(x, y)|^2 d^\nu x d^\nu y \leq \|f\|_{L^2(\mathbb{R}^\nu)}^2 \|\chi_{\Omega_m}\|_{L^2(\mathbb{R}^\nu)} < \infty$$

since  $|\Omega_m| < \infty$ . Thus,  $C_m$  is Hilbert–Schmidt, so compact.  $\square$

*Proof of Theorem 2.* We can follow the proof of Theorem 1. It suffices to prove that  $e^{-sL_0} e^{-sV}$  is compact, and so, that  $e^{-sL_0} \chi_{\Omega_m}$  is Hilbert–Schmidt.

That  $e^{-sL_0}$  maps  $L^2$  to  $L^\infty$  implies, according to the Dunford–Pettis theorem (see [24, Thm. 46.1]), that there is, for each  $x \in X$ , a function  $f_x(\cdot) \in L^2(X, d\mu)$  with

$$(3.6) \quad (e^{-sL_0} g)(x) = \langle f_x, g \rangle$$

and

$$(3.7) \quad \sup_x \|f_x\|_{L^2} = \|e^{-sL_0}\|_{L^2 \rightarrow L^\infty} \equiv C < \infty.$$

Thus,  $e^{-sL_0}$  has an integral kernel  $K(x, y)$  with

$$(3.8) \quad \sup_x \int |K(x, y)|^2 d\mu(y) = C < \infty$$

(for  $K(x, y) = f_x(y)$ ). But  $e^{-sL_0}$  is selfadjoint, so its kernel is complex symmetric, so

$$(3.9) \quad \sup_y \int |K(x, y)|^2 d\mu(x) = C < \infty.$$

Thus,

$$(3.10) \quad \int |K(x, y)\chi_{\Omega_m}(y)|^2 d\mu(x)d\mu(y) \leq C\mu(\Omega_m) < \infty$$

and  $e^{-sL_0}\chi_{\Omega_m}$  is Hilbert–Schmidt.  $\square$

#### 4. PROOF OF THEOREM 3

As with the proof of Theorem 1, it suffices to prove that for each  $M$ ,  $e^\Delta \chi_{\Omega_M}$  is compact.  $e^\Delta$  is convolution with an  $L^1$  function,  $f$ . Let  $Q_R$  be the characteristic function of  $\{x \mid |x| < R\}$ . Let  $F_R$  be convolution with  $fQ_R$ . Then

$$(4.1) \quad \|e^\Delta - F_R\| \leq \|f(1 - Q_R)\|_1 \rightarrow 0$$

as  $R \rightarrow \infty$ , so

$$(4.2) \quad \|e^\Delta \chi_{\Omega_M} - F_R \chi_{\Omega_M}\| \rightarrow 0$$

and it suffices to prove for each  $R, M$ ,

$$(4.3) \quad C_{M,R} = F_R \chi_{\Omega_M}$$

is compact. Clearly, this works if we show for some  $k$ ,  $(C_{M,R}^* C_{M,R})^k$  is Hilbert–Schmidt.

Let  $D$  be the operator with integral kernel

$$(4.4) \quad D(x, y) = \chi_{\Omega_M}(x) Q_{2R}(x - y) \chi_{\Omega_M}(y).$$

Since  $f$  is bounded, it is easy to see that

$$(4.5) \quad (C_{M,R}^* C_{M,R})(x, y) \leq c D(x, y)$$

for some constant  $c$ , so it suffices to show  $D^k$  is Hilbert–Schmidt.

$D^k$  has integral kernel

$$(4.6) \quad D^k(x, y) = \int D(x, x_1) D(x_1, x_2) \dots D(x_{k-1}, y) dx_1 \dots dx_{k-1}.$$

Fix  $y$ . This integral is zero unless  $|x - x_1| < 2R, \dots, |x_{k-1} - y| < 2R$ , so, in particular, unless  $|x - y| \leq 2kR$ . Moreover, the integrand can certainly be restricted to the regions  $|x_j - y| \leq 2kR$ . Thus,

$$(4.7) \quad D^k(x, y) \leq Q_{2kR}(x - y) \left( \int_{|x_j - y| \leq 2kR} \prod_{j=1}^{k-1} \chi_{\Omega_M}(x_j) dx_1 \dots dx_{k-1} \right) \chi_{\Omega_m}(y)$$

$$(4.8) \quad = Q_{2kR}(x - y) (\omega_y^{2kR} (\Omega_M)^{k-1}) \chi_{\Omega_m}(y)$$

by the definition of  $\omega_x^\ell$  in (1.7).

Thus,

$$\int |D^k(x, y)|^2 d^\nu x d^\nu y \leq C(kR)^\nu \int_{x \in \Omega} [\omega_x^{2kR} (\Omega_M)]^{2k-2} d^\nu x$$

so if  $2k - 2 > r$  and (1.9) holds,  $D^k$  is Hilbert–Schmidt.  $\square$

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