# ON THE COMMON POINT SPECTRUM OF PAIRS OF SELF-ADJOINT EXTENSIONS

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Dedicated to Vladimir Koshmanenko on the occasion of his 70th birthday

ABSTRACT. Given two different self-adjoint extensions of the same symmetric operator, we analyse the intersection of their point spectra. Some simple examples are provided.

#### 1. Preliminaries

Given a linear closed operator L, we denote by

$$\mathcal{D}(L)$$
,  $\mathcal{K}(L)$ ,  $\mathcal{R}(L)$ ,  $\mathcal{G}(L)$ ,  $\rho(L)$ 

its domain, kernel, range, graph and resolvent set respectively.  $\mathcal{H}$  denotes a Hilbert space with scalar product  $\langle \cdot, \cdot \rangle$  and corresponding norm  $\| \cdot \|$ ; we also make use of an auxiliary Hilbert space  $\mathfrak{h}$  with scalar product  $(\cdot, \cdot)$  and corresponding norm  $| \cdot |$ .

Given a closed, densely defined, symmetric operator

$$S: \mathcal{D}(S) \subseteq \mathcal{H} \to \mathcal{H}$$

with equal deficiency indices, by von Neumann's theory one has (here the direct sums are given w.r.t. the graph inner product of  $S^*$ )

$$\mathcal{D}(S^*) = \mathcal{D}(S) \oplus \mathcal{K}_+ \oplus \mathcal{K}_-, \quad \mathcal{K}_{\pm} := \mathcal{K}(-S^* \pm i),$$

$$S^*(\phi_{\circ} \oplus \phi_{+} \oplus \phi_{-}) = S\phi_{\circ} + i\phi_{+} - i\phi_{-},$$

and any self-adjoint extension of S is of the kind  $A_U = S^* | \mathcal{G}(U)$ , the restriction of  $S^*$  to  $\mathcal{G}(U)$ , where  $U : \mathcal{K}_+ \to \mathcal{K}_-$  is unitary. Therefore, fixing a unitary  $U_\circ$  and posing  $A := A_{U_\circ}$ , one has

$$S = A | \mathcal{K}(\tau_{\circ}), \quad \tau_{\circ} : \mathcal{D}(A) \to \mathfrak{h}_{\circ},$$

where

$$\mathfrak{h}_{\circ} = \mathcal{K}_{+} \,, \quad \tau_{\circ} = P_{+} \,,$$

and  $P_+$  is the orthogonal (w.r.t. the graph inner product of  $S^*$ ) projection onto  $\mathcal{K}_+$ . Since  $\mathcal{K}(\tau_\circ) = \mathcal{K}(\tau)$  where  $\tau = M\tau_\circ$  and  $M : \mathfrak{h}_\circ \to \mathfrak{h}$  is any continuous linear bijection, in the search of the self-adjoint extension of S, we can consider the following equivalent problem: determine all the self-adjoint extensions of  $A|\mathcal{K}(\tau)$ , where

$$\tau: \mathcal{D}(A) \to \mathfrak{h}$$

is a linear, continuous (with respect to the graph norm on  $\mathcal{D}(A)$ ), surjective map onto an auxiliary Hilbert space  $\mathfrak{h}$  with its kernel  $\mathcal{K}(\tau)$  dense in  $\mathcal{H}$ . Typically A is a differential operator,  $\tau$  is some trace (restriction) operator along a null subset N and  $\mathfrak{h}$  is some function space over N.

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We suppose that the spectrum of A does not coincide with the whole real line and so, by eventually adding a constant to A, we make the following hypothesis:

$$0 \in \rho(A)$$
.

By the results provided in [9] and [7] (to which we refer for proofs and connections with equivalent formulations, in particular with boundary triplets theory) one has the following

**Theorem 1.1.** The set of all self-adjoint extensions of S is parametrized by the set  $E(\mathfrak{h})$ of couples  $(\Pi, \Theta)$ , where  $\Pi$  is an orthogonal projection in  $\mathfrak{h}$  and  $\Theta$  is a self-adjoint operator in  $\mathcal{R}(\Pi)$ . If  $A^{\Pi,\Theta}$  denotes the self-adjoint extension corresponding to  $(\Pi,\Theta) \in \mathsf{E}(\mathfrak{h})$  then

$$A^{\Pi,\Theta}: \mathcal{D}(A^{\Pi,\Theta}) \subseteq \mathcal{H} \to \mathcal{H}, \quad A^{\Pi,\Theta}\phi := A\phi_0$$

$$\mathcal{D}(A^{\Pi,\Theta}) := \{ \phi = \phi_0 + G_0 \xi_{\phi}, \ \phi_0 \in \mathcal{D}(A), \xi_{\phi} \in \mathcal{D}(\Theta), \ \Pi \tau \phi_0 = \Theta \xi_{\phi} \} ,$$

where

$$G_z: \mathfrak{h} \to \mathcal{H}, \quad G_z:= \left(\tau(-A+\bar{z})^{-1}\right)^*, \quad z \in \rho(A).$$

Moreover the resolvent of  $A^{\Pi,\Theta}$  is given, for any  $z \in \rho(A) \cap \rho(A^{\Pi,\Theta})$ , by the Krein's type formula

$$(-A^{\Pi,\Theta}+z)^{-1} = (-A+z)^{-1} + G_z \Pi(\Theta + z \Pi G_0^* G_z \Pi)^{-1} \Pi G_{\bar{z}}^* \,.$$

**Remark 1.2.** Notice that the extension corresponding to  $\Pi = 0$  is A itself. The extension corresponding to  $(1,\Theta)$  is denoted by  $A^{\Theta}$  and everywhere we omit the index  $\Pi$  in the case  $\Pi = 1$ . By [8], Corollary 3.2, the sub-family  $\{A^{\Theta} : \Theta \text{ self-adjoint}\}\$  gives all singular perturbations of A, where we say that  $\hat{A}$  is a singular perturbation of A whenever the set  $\{\phi \in \mathcal{D}(A) \cap \mathcal{D}(\hat{A}) : A\phi = \hat{A}\phi\}$  is dense in  $\mathcal{H}$  (see [4]).

**Remark 1.3.** The operator  $G_z$  is injective (by surjectivity of  $\tau$ ) and for any  $z \in \rho(A)$ one has (see [7], Remark 2.8)

$$\mathcal{R}(G_z) \cap \mathcal{D}(A) = \{0\},\,$$

so that the decomposition appearing in  $\mathcal{D}(A^{\Pi,\Theta})$  is unique. Moreover (see [7], Lemma 2.1)

(1.2) 
$$G_w - G_z = (z - w)(-A + w)^{-1}G_z.$$

## 2. The common point spectrum

Given a self-adjoint operator A let us denote by

$$\sigma(A)$$
,  $\sigma_p(A)$ ,  $\sigma_d(A)$ 

its full, point and discrete spectrum respectively.

Given  $\lambda \in \sigma_p(A)$ , we denote by  $P_{\lambda}$  the orthogonal projector onto the corresponding eigenspace  $\mathcal{H}_{\lambda} \subseteq \mathcal{D}(A)$  and pose  $P_{\lambda}^{\perp} := 1 - P_{\lambda}$ . Given  $\lambda \in \sigma_p(A^{\Pi,\Theta})$ , we denote by  $\mathcal{H}_{\lambda}^{\Pi,\Theta} \subseteq \mathcal{D}(A^{\Pi,\Theta})$  the corresponding eigenspace. As regards the eigenvalues of  $A^{\Pi,\Theta}$  which are not in the spectrum of A a complete

answer is given by the following result which is consequence of Kreĭn's resolvent formula (see [3], Section 2, Propositions 1 and 2, and [8], Theorem 3.4):

## Lemma 2.1.

$$\lambda \in \rho(A) \cap \sigma_p(A^{\Pi,\Theta}) \iff 0 \in \sigma_p(\Theta + \lambda \Pi G_0^* G_\lambda \Pi),$$
$$\mathcal{H}_{\lambda}^{\Pi,\Theta} = \{ G_{\lambda} \xi \,, \, \xi \in \mathcal{K}(\Theta + \lambda \Pi G_0^* G_\lambda \Pi) \}.$$

Here we are interested in the common eigenvalues, i.e. in the points in  $\sigma_p(A) \cap \sigma_p(A^{\Pi,\Theta})$ . Therefore we take  $\lambda \in \sigma_p(A)$  and we look for solutions  $\phi \in \mathcal{D}(A^{\Pi,\Theta})$  of the eigenvalue equation

$$A^{\Pi,\Theta}\phi = \lambda\phi$$
,

i.e., by Theorem 1.1,

$$(A - \lambda)\phi_0 = \lambda G_0 \xi_\phi$$
.

By

$$(A - \lambda)P_{\lambda}\phi_0 = 0$$
,  $(A - \lambda)P_{\lambda}^{\perp}\phi_0 \in \mathcal{R}(P_{\lambda}^{\perp})$ ,

this is equivalent to the couple of equations

$$(2.1) P_{\lambda} G_0 \xi_{\phi} = 0,$$

$$(2.2) (A - \lambda)P_{\lambda}^{\perp}\phi_0 = \lambda P_{\lambda}^{\perp}G_0\xi_{\phi}$$

together with the constraint

(2.3) 
$$\xi_{\phi} \in \mathcal{D}(\Theta) \subseteq \mathcal{R}(\Pi), \quad \Pi \tau \phi_0 = \Theta \xi_{\phi}.$$

Equation (2.1) gives, for all  $\psi \in \mathcal{H}$ ,

$$0 = \langle G_0 \xi_{\phi}, P_{\lambda} \psi \rangle = -\langle \xi_{\phi}, \tau A^{-1} P_{\lambda} \psi \rangle = -\frac{1}{\lambda} \langle \xi_{\phi}, \tau P_{\lambda} \psi \rangle$$

and so

$$\xi_{\phi} \in (\mathcal{R}(\tau P_{\lambda}))^{\perp}$$
.

If  $\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp} = \{0\}$  then, since  $G_0$  is injective, one has that in this case  $\phi$  is an eigenvector with eigenvalue  $\lambda$  if and only if  $\phi \in \mathcal{H}_{\lambda}$  and  $\Pi \tau \phi = 0$ .

Conversely suppose that  $\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp} \neq \{0\}$  and moreover that  $\lambda$  is an isolated eigenvalue. Then  $\lambda \in \rho(A|\mathcal{H}_{\lambda}^{\perp})$  and (2.2) gives

$$P_{\lambda}^{\perp}\phi_0 = -\lambda(-A+\lambda)^{-1}P_{\lambda}^{\perp}G_0\xi_{\phi}.$$

By  $\Pi \tau \phi_0 = \Theta \xi_{\phi}$  then one gets

$$(2.4) \qquad \Pi \tau P_{\lambda} \phi_0 = \Theta \xi_{\phi} - \Pi \tau P_{\lambda}^{\perp} \phi_0 = (\Theta + \lambda \Pi \tau (-A + \lambda)^{-1} P_{\lambda}^{\perp} G_0 \Pi) \xi.$$

By defining

$$G_{\lambda}^{\perp}: \mathfrak{h} \to \mathcal{H} \,, \quad G_{\lambda}^{\perp}:= (\tau(-A+\lambda)^{-1}P_{\lambda}^{\perp})^* \,,$$

and by  $(G_{\lambda}^{\perp})^*G_0 = G_0^*G_{\lambda}^{\perp}$  (this relation is consequence of (1.2)), (2.4) is equivalent to

$$\Pi \tau P_{\lambda} \phi_0 = (\Theta + \lambda \Pi G_0^* G_{\lambda}^{\perp} \Pi) \xi$$
.

Moreover by  $(-A + \lambda)^{-1} P_{\lambda}^{\perp} G_0 = -A^{-1} G_{\lambda}^{\perp}$  one has

$$P_{\lambda}\phi_{0} + P_{\lambda}^{\perp}\phi_{0} + G_{0}\xi_{\phi} = P_{\lambda}\phi_{0} + (-\lambda(-A+\lambda)^{-1}P_{\lambda}^{\perp} + P_{\lambda}^{\perp})G_{0}\xi_{\phi}$$
$$= P_{\lambda}\phi_{0} - A(-A+\lambda)^{-1}P_{\lambda}^{\perp}G_{0}\xi_{\phi}$$
$$= P_{\lambda}\phi_{0} + G_{\lambda}^{\perp}\xi_{\phi}.$$

In conclusion we have proven the following

Theorem 2.2. Let  $\lambda \in \sigma_p(A)$ .

1) Suppose

$$\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp} = \{0\}$$

and pose

$$\mathcal{K}_{\lambda}^{\Pi} := \{ \psi \in \mathcal{H}_{\lambda} : \Pi \tau \psi = 0 \}.$$

Then

$$\lambda \in \sigma_p(A^{\Pi,\Theta}) \quad \iff \quad \mathcal{K}_{\lambda}^{\Pi} \neq \{0\}$$

and

$$\mathcal{H}_{\lambda}^{\Pi,\Theta} = \mathcal{K}_{\lambda}^{\Pi}$$
.

2) Suppose

$$\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp} \neq \{0\}$$

and let  $\lambda$  be isolated. Let  $\mathcal{N}_{\lambda}^{\Pi,\Theta}$  be the set of couples  $(\psi,\xi) \in \mathcal{H}_{\lambda} \oplus \mathcal{R}(\Pi)$  such that

(2.5) 
$$\xi \in D(\Theta) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp},$$

(2.6) 
$$\Pi \tau \psi = (\Theta + \lambda \Pi G_0^* G_\lambda^{\perp} \Pi) \xi.$$

Then

$$\lambda \in \sigma_p(A^{\Pi,\Theta}) \iff \mathcal{N}_{\lambda}^{\Pi,\Theta} \neq \{0\},$$
  
$$\dim(\mathcal{H}_{\lambda}^{\Pi,\Theta}) = \dim(\mathcal{N}_{\lambda}^{\Pi,\Theta})$$

and

$$\mathcal{H}_{\lambda}^{\Pi,\Theta} = \{ \phi \in \mathcal{H} : \phi = \psi + G_{\lambda}^{\perp} \xi, \ (\psi, \xi) \in \mathcal{N}_{\lambda}^{\Pi,\Theta} \}.$$

Remark 2.3. Notice that

$$(\mathcal{R}(\Pi))^{\perp} \cap \mathcal{R}(\tau P_{\lambda}) \neq \{0\} \implies \mathcal{K}_{\lambda}^{\Pi} \neq \{0\}.$$

**Remark 2.4.** Suppose  $\lambda \in \sigma_n(A)$  is isolated. Noticing that

$$\mathcal{K}_{\lambda}^{\Pi} \oplus (\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp} \cap \mathcal{K}(\Theta + \lambda \Pi G_{0}^{*} G_{\lambda}^{\perp} \Pi)) \subseteq \mathcal{N}_{\lambda}^{\Pi,\Theta},$$

one has

$$\mathcal{K}_{\lambda}^{\Pi} \neq \{0\} \implies \lambda \in \sigma_p(A^{\Pi,\Theta}).$$

In particular, in the case  $\lambda$  is simple with eigenvector  $\psi_{\lambda}$ ,

$$\Pi \tau \psi_{\lambda} = 0 \implies \lambda \in \sigma_p(A^{\Pi,\Theta}).$$

**Remark 2.5.** Suppose  $\mathcal{R}(\tau P_{\lambda}) = \{0\}$ . Then  $\mathcal{K}_{\lambda}^{\Pi} = \mathcal{H}_{\lambda}$  and so, in case  $\lambda \in \sigma_p(A)$  is isolated,  $\lambda \in \sigma_p(A^{\Pi,\Theta})$  and

$$\mathcal{H}_{\lambda}^{\Pi,\Theta} = \left\{ \phi = \psi_{\lambda} + G_{\lambda}^{\perp} \xi \,, \ \psi_{\lambda} \in \mathcal{H}_{\lambda} \,, \ \xi \in \mathcal{R}(\Pi) \cap \mathcal{K}(\Theta + \lambda \Pi G_0^* G_{\lambda}^{\perp} \Pi) \right\}.$$

Remark 2.6. The papers [1] and [6] contain results related to the ones given by Theorem 2.2 (see Theorem 3.6 in [6] and Theorem 4.7 in [1]). We thank Konstantin Pankrashkin for the communication.

# 3. Examples

3.1. Rank-one singular perturbations. Suppose  $\mathfrak{h} = \mathbb{C}$ . Then  $\Pi = 1, \Theta = \theta \in \mathbb{R}$ and either  $\mathcal{R}(\tau P_{\lambda}) = \mathbb{C}$  or  $\mathcal{R}(\tau P_{\lambda}) = \{0\}.$ 

If  $\mathcal{R}(\tau P_{\lambda}) = \mathbb{C}$  then  $\lambda \in \sigma_p(A^{\theta})$  if and only if  $\mathcal{K}_{\lambda} \neq \{0\}$ , where

$$\mathcal{K}_{\lambda} := \{ \psi \in \mathcal{H}_{\lambda} : \tau \psi = 0 \}.$$

Since  $\mathcal{R}(\tau P_{\lambda}) = \{0\}$  if and only if  $\mathcal{K}_{\lambda} = \mathcal{H}_{\lambda}$ , when  $\lambda$  is isolated and  $\mathcal{R}(\tau P_{\lambda}) = \{0\}$  one

$$\mathcal{N}_{\lambda}^{\theta} = \mathcal{K}_{\lambda} \oplus \{ \xi \in \mathbb{C} : (\theta + \lambda \langle G_0, G_{\lambda}^{\perp} \rangle) \xi = 0 \}$$

and so

$$\theta + \lambda \langle G_0, G_{\lambda}^{\perp} \rangle = 0 \quad \Longrightarrow \quad \mathcal{N}_{\lambda}^{\theta} = \mathcal{K}_{\lambda} \oplus \mathbb{C} \equiv \mathcal{H}_{\lambda} \oplus \mathbb{C},$$
$$\theta + \lambda \langle G_0, G_{\lambda}^{\perp} \rangle \neq 0 \quad \Longrightarrow \quad \mathcal{N}_{\lambda}^{\theta} = \mathcal{K}_{\lambda} \oplus \{0\} \equiv \mathcal{H}_{\lambda}.$$

In conclusion when  $\mathfrak{h} = \mathbb{C}$  and  $\lambda \in \sigma_p(A)$  is isolated,

$$(3.1) \lambda \in \sigma_p(A^{\theta}) \iff \mathcal{K}_{\lambda} \neq \{0\}$$

and

$$\mathcal{H}^{\theta}_{\lambda} = \{ \psi = \psi_{\lambda} + G^{\perp}_{\lambda} \xi, \ \psi_{\lambda} \in \mathcal{H}_{\lambda}, \ (\theta + \lambda \langle G_0, G^{\perp}_{\lambda} \rangle) \xi = 0 \}.$$

In particular if  $\lambda$  is a simple isolated eigenvalue of A with corresponding eigenfunction  $\psi_{\lambda}$ , then  $\lambda \in \sigma_p(A^{\theta})$  if and only if  $\tau\psi_{\lambda} = 0$ . For example, if  $\mathcal{H} = L^2(\Omega)$  and  $\tau : \mathcal{D}(A) \to \mathbb{C}$  is the evaluation map at  $y \in \Omega$ ,  $\tau\psi := \psi(y)$ , then  $\lambda$  is preserved if and only if y belongs to the nodal set (if any) of  $\psi_{\lambda}$ . Thus if A is (minus) the Dirichlet Laplacian on a bounded open set  $\Omega \subset \mathbb{R}^d$ ,  $d \leq 3$ , its lowest eigenvalue is never preserved under a point perturbation. Analogous results hold in the case A is the Laplace-Beltrami operator on a compact d-dimensional Riemannian manifold M,  $d \leq 3$ , thus reproducing the ones given in [2], Theoreme 2, part 1.

## 3.2. The Šeba billiard. Let

$$A = \Delta : \mathcal{D}(A) \subset L^2(R) \to L^2(R)$$
,

$$\mathcal{D}(A) = \{ \phi \in C(\overline{R}) : \Delta \phi \in L^2(R), \ \phi(\mathsf{x}) = 0, \ \mathsf{x} \in \partial R \},$$

be the Dirichlet Laplacian on the rectangle  $R = (0, a) \times (0, b)$ . Then

$$\sigma(A) = \sigma_d(A) = \{\lambda_{m,n}, (m,n) \in \mathbb{N}^2\}$$

and

$$\mathcal{H}_{\lambda_{m,n}} = \operatorname{span}\{\psi_{m',n'} : \lambda_{m',n'} = \lambda_{m,n}\},\,$$

where

$$\lambda_{m,n} := -\pi^2 \left( \frac{m^2}{a^2} + \frac{n^2}{b^2} \right)$$

and

$$\psi_{m,n}(\mathsf{x}) := \sin\left(\frac{m\pi x_1}{a}\right) \sin\left(\frac{n\pi x_2}{b}\right), \quad \mathsf{x} \equiv (x_1, x_2).$$

Let

$$\tau\psi:=\psi(\mathsf{y})\,,$$

so that  $A^{\theta}$  describes a "Šeba billiard", i.e. the Dirichlet Laplacian on the rectangle R with a point perturbation placed at the point  $y \equiv (y_1, y_2)$  (see [10]).

Since  $\sigma(A) = \sigma_d(A)$ , by the invariance of the essential spectrum under finite rank perturbations, one has  $\sigma(A^{\theta}) = \sigma_d(A^{\theta})$  and, by (3.1),  $\lambda_{m,n} \in \sigma(A) \cap \sigma(A^{\theta})$  if and only

$$\forall (m', n') \text{ s.t. } \lambda_{m', n'} = \lambda_{m, n}, \quad \sin\left(\frac{m'\pi y_1}{a}\right) \sin\left(\frac{n'\pi y_2}{b}\right) = 0.$$

Equivalently

$$\sigma(A) \cap \sigma(A^{\theta}) = \emptyset \quad \iff \quad \left(\frac{y_1}{a}, \frac{y_2}{b}\right) \notin \mathbb{Q}^2.$$

If there exists relatively prime integers  $1 \le p < q$  such that  $\frac{y_1}{a} = \frac{p}{q}$  while  $\frac{y_2}{b}$  is irrational, then

$$\sigma(A) \cap \sigma(A^{\theta}) = \{\lambda_{kq,n}, (k,n) \in \mathbb{N}^2\}.$$

Analogously if  $\frac{y_1}{a}$  is irrational and  $\frac{y_2}{b} = \frac{p}{a}$  then

$$\sigma(A) \cap \sigma(A_{\theta}) = \{\lambda_{m,kq}, (m,k) \in \mathbb{N}^2\}$$

while if  $\frac{y_1}{a} = \frac{p}{a}$  and  $\frac{y_2}{b} = \frac{r}{s}$ , then

$$\sigma(A) \cap \sigma(A^{\theta}) = \{\lambda_{ka,n}, (k,n) \in \mathbb{N}^2\} \cup \{\lambda_{m,ks}, (m,k) \in \mathbb{N}^2\}.$$

- 3.3. Rank-two singular perturbations. Let  $\mathfrak{h} = \mathbb{C}^2$ . Then either  $\Pi = 1$  or  $\Pi = w \otimes w$ ,  $w \in \mathbb{C}^2$ , |w| = 1. Let  $\lambda \in \sigma_p(A)$ .
- 1.1)  $\mathcal{R}(\tau P_{\lambda}) = \mathbb{C}^2$ ,  $\Pi = 1$ . Then  $\lambda \in \sigma_p(A^{\Theta})$  if and only if there exists  $\psi \in \mathcal{H}_{\lambda} \setminus \{0\}$  such that  $\tau \psi = 0$ .
- 1.2)  $\mathcal{R}(\tau P_{\lambda}) = \mathbb{C}^2$ ,  $\Pi = w \otimes w$ . Then  $\lambda \in \sigma_p(A^{\Pi,\Theta})$  if and only if there exists  $\psi \in \mathcal{H}_{\lambda} \setminus \{0\}$  such that  $w \cdot \tau \psi = 0$ .

Now suppose further that  $\lambda \in \sigma_p(A)$  is isolated.

2.1)  $\mathcal{R}(\tau P_{\lambda}) = \operatorname{span}(\xi_{\lambda}) \simeq \mathbb{C}, \ |\xi_{\lambda}| = 1, \ \Pi = 1.$  Decomposing equation (2.6) w.r.t. the orthonormal base  $\{\xi_{\lambda}, \xi_{\lambda}^{\perp}\}$  one gets that  $\mathcal{N}_{\lambda}^{\Theta} \neq \{0\}$  if and only if there exists  $\zeta \equiv (\zeta_{1}, \zeta_{2}) \in \mathbb{C}^{2} \setminus \{0\}$  solving

$$\begin{cases} \zeta_1 = (\xi_\lambda \cdot (\Theta + \lambda G_0^* G_\lambda^\perp) \xi_\lambda^\perp) \zeta_2, \\ 0 = (\xi_\lambda^\perp \cdot (\Theta + \lambda G_0^* G_\lambda^\perp) \xi_\lambda^\perp) \zeta_2 \,. \end{cases}$$

Hence

$$\lambda \in \sigma_p(A^\Theta) \iff (\xi_\lambda^\perp \cdot (\Theta + \lambda G_0^* G_\lambda^\perp) \xi_\lambda^\perp) = 0 \,.$$

2.2)  $\mathcal{R}(\tau P_{\lambda}) = \operatorname{span}(\xi_{\lambda}) \simeq \mathbb{C}$ ,  $\Pi = w \otimes w$ . Let us use the decomposition  $w = w_{||} + w_{\perp}$  w.r.t. the orthonormal base  $\{\xi_{\lambda}, \xi_{\lambda}^{\perp}\}$ . If  $w_{||} = 0$  then  $\mathcal{K}_{\lambda}^{\Pi} \neq \{0\}$  and so  $\lambda \in \sigma_{p}(A^{\Pi,\Theta})$ . If  $w_{||} \neq 0$  then  $\mathcal{K}_{\lambda}^{\Pi} = \{0\}$  and  $\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda}))^{\perp} = \{0\}$ , thus  $\lambda \notin \sigma_{p}(A^{\Pi,\Theta})$ . In conclusion

$$\lambda \in \sigma_p(A^{\Pi,\Theta}) \quad \iff \quad w = \xi_\lambda^\perp.$$

- 3)  $\mathcal{R}(\tau P_{\lambda}) = \{0\}$ . In this case  $\lambda \in \sigma_p(A^{\Pi,\Theta})$ .
- 3.4. The Laplacian on a bounded interval. Let

$$A: \mathcal{D}(A) \subseteq L^2(0,a) \to L^2(0,a), \quad A\phi = \phi'',$$

$$\mathcal{D}(A) = \{ \phi \in C^1[0, a] : \phi'' \in L^2(0, a), \ \phi(0) = \phi(a) = 0 \},\$$

be the Dirichlet Laplacian on the bounded interval (0, a) and pose

$$\tau: \mathcal{D}(A) \to \mathbb{C}^2$$
,  $\tau \phi \equiv \gamma_1 \phi := (\phi'(0), -\phi'(a))$ .

Therefore  $S = A | \mathcal{K}(\tau)$  is the minimal Laplacian with domain

$$\mathcal{D}(S) = \{ \phi \in C^1[0, a] : \phi'' \in L^2(0, a), \ \phi(0) = \phi'(0) = \phi(a) = \phi'(a) = 0 \}$$

and the self-adjoint extensions of S are rank-two perturbations of the Dirichlet Laplacian A. One has

$$\sigma(A) = \sigma_d(A) = \{\lambda_n\}_1^{\infty}, \quad \lambda_n = -\left(\frac{n\pi}{a}\right)^2$$

and the normalized eigenvector corresponding to  $\lambda_n$  is

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right).$$

By Theorem 1.1 and by the change of extension parameter (here  $P_0$  represents the Dirichlet-to-Neumann operator)

$$(\Pi, \Theta) \mapsto (\Pi, B), \quad B := \Theta - \Pi P_0 \Pi, \quad P_0 \equiv \frac{1}{a} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}$$

any self-adjoint extension of the minimal Laplacian S is of the kind  $A^{\Pi,B}$ ,  $(\Pi,B) \in \mathsf{E}(\mathbb{C}^2)$ , where

$$A^{\Pi,B}:\mathcal{D}(A^{\Pi,B})\subset L^2(0,a)\to L^2(0,a)\,,\quad A^{\Pi,B}\phi=\phi''\,,$$

$$\mathcal{D}(A^{\Pi,B}) = \{ \phi \in C^1[0,a] : \phi'' \in L^2(0,a) , \ \gamma_0 \phi \in \mathcal{R}(\Pi) , \ \Pi \gamma_1 \phi = B \gamma_0 \phi \}$$

(see e.g. [9], Example 5.1). Here  $\gamma_0 \phi := (\phi(0), \phi(a))$ .

The case  $\Pi=0$  reproduces A itself, the case  $\Pi=1,\ B=\begin{pmatrix}b_{11}&b_{12}\\\bar{b}_{12}&b_{22}\end{pmatrix},\ b_{11},b_{22}\in\mathbb{R},$   $b_{12}\in\mathbb{C}$ , gives the boundary conditions

$$\begin{cases} b_{11} \phi(0) - \phi'(0) + b_{12} \phi(a) = 0, \\ \bar{b}_{12} \phi(0) + b_{22} \phi(a) + \phi'(a) = 0, \end{cases}$$

and the case  $\Pi = w \otimes w$ ,  $w \equiv (w_1, w_2) \in \mathbb{C}^2$ ,  $|w_1|^2 + |w_2|^2 = 1$ ,  $B \equiv b \in \mathbb{R}$ , gives the boundary conditions

$$\begin{cases} w_2 \phi(0) - w_1 \phi(a) = 0, \\ \bar{w}_1 (b \phi(0) - \phi'(0)) + \bar{w}_2 (b \phi(a) + \phi'(a)) = 0. \end{cases}$$

By the invariance of the essential spectrum under finite rank perturbations,  $\sigma(A^{\Pi,B}) = \sigma_d(A^{\Pi,B})$ . Now we use the results given in subsection 3.3. One has

$$\mathcal{R}(\tau P_{\lambda_n}) = \operatorname{span}(\hat{\xi}_n), \quad \hat{\xi}_n \equiv \frac{1}{\sqrt{2}} \left( 1, (-1)^{n-1} \right).$$

Let  $\Pi=1$  and  $\hat{\xi}_n^{\perp}\equiv\frac{1}{\sqrt{2}}(1,(-1)^n)$ . By point 2.1 in subsection 3.3 we known that  $\lambda_n\in\sigma(A^B)$  if and only if  $\hat{\xi}_n^{\perp}\cdot(B+P_0+\lambda_nG_0^*G_{\lambda_n}^{\perp})\hat{\xi}_n^{\perp}=0$ . Since the resolvent of A is explicitly known,  $\hat{\xi}_n^{\perp}\cdot(B+P_0+\lambda_nG_0^*G_{\lambda_n}^{\perp})\hat{\xi}_n^{\perp}$  can be calculated. However we use here a short cut which avoids any calculation: the Neumann Laplacian corresponds to B=0 and we know that its spectrum is  $\{0\}\cup\sigma(A)$ , thus

$$\hat{\xi}_n^{\perp} \cdot (P_0 + \lambda_n G_0^* G_{\lambda_n}^{\perp}) \hat{\xi}_n^{\perp} = 0.$$

Therefore we obtain

$$\lambda_n \in \sigma(A^B) \iff b_{11} + b_{22} + 2(-1)^n \text{Re}(b_{12}) = 0.$$

If  $\Pi = w \otimes w$  by point 2.2 in subsection 3.3 one has

$$\lambda_n \in \sigma(A^{\Pi,B}) \quad \iff \quad w = \hat{\xi}_n^{\perp}.$$

In both cases

$$\lambda_n \in \sigma(A^{\Pi,B}) \iff \lambda_{n+2} \in \sigma(A^{\Pi,B}).$$

Moreover

$$\sigma(A) \subseteq \sigma(A^{\Pi,B}) \iff \Pi = 1 \text{ and } b_{11} + b_{22} = 0, \operatorname{Re}(b_{12}) = 0.$$

3.5. Equilateral quantum graphs. Let  $\mathcal{H} = \bigoplus_{k=1}^{N} L^2(0, a)$  and  $A_N = \bigoplus_{k=1}^{N} A$ , where A is defined as in subsection 3.4 (to which we refer for notations). Then  $\sigma(A_N) = \sigma_d(A_N) = \sigma(A)$  and the eigenfunctions corresponding to the N-fold degenerate eigenvalue  $\lambda_n$  are

$$\Psi_{k,n} = \bigoplus_{i=1}^{N} \psi_{i,k,n}, \quad k = 1, \dots, N, \quad \psi_{i,k,n} = \begin{cases} 0, & i \neq k, \\ \psi_n, & i = k. \end{cases}$$

By taking

$$\tau: \mathcal{D}(A_N) \equiv \oplus_{k=1}^N \mathcal{D}(A) \to \oplus_{k=1}^N \mathbb{C}^2 \equiv \mathbb{C}^{2N} \,, \quad \tau = \oplus_{k=1}^N \gamma_1 \,,$$

one gets, by Theorem 1.1, self-adjoint extensions describing quantum graphs (see e.g. [5]) with N edges of the same length a. By Theorem 1.1 and by the change of extension parameter

$$(\Pi, \Theta) \mapsto (\Pi, B), \quad B := \Theta - \Pi(\bigoplus_{k=1}^{N} P_0)\Pi,$$

such extensions are of the kind  $A^{\Pi,B}$ ,  $(\Pi,B) \in \mathsf{E}(\mathbb{C}^{2N})$ , where (see [9], Example 5.2).

$$A^{\Pi,B}: \mathcal{D}(A^{\Pi,B}) \subset \bigoplus_{k=1}^{N} L^{2}(0,a) \to \bigoplus_{k=1}^{N} L^{2}(0,a),$$
$$A^{\Pi,B}(\bigoplus_{k=1}^{N} \phi_{k}) = \bigoplus_{k=1}^{N} \phi_{k}'',$$

$$\mathcal{D}(A^{\Pi,B}) = \left\{ \bigoplus_{k=1}^{N} \phi_k : \phi_k \in C^1[0,a], \ \phi_k'' \in L^2(0,a), \right. \\ \left. \left( \bigoplus_{k=1}^{N} \gamma_0 \phi_k \right) \in \mathcal{R}(\Pi), \ \Pi(\bigoplus_{k=1}^{N} \gamma_1 \phi_k) = B(\bigoplus_{k=1}^{N} \gamma_0 \phi_k) \right\}.$$

The couple  $(\Pi, B)$  represents the connectivity of the quantum graph.

1)  $\Pi = 1$ . Given  $\lambda_n \in \mathcal{D}(A)$ , we pose

$$\mathbb{C}^{2N}_{||} := \oplus_{k=1}^N \mathrm{span}(\hat{\xi}_n) \simeq \mathbb{C}^N \,, \quad \mathbb{C}^{2N}_{\perp} := \oplus_{k=1}^N \mathrm{span}(\hat{\xi}_n^{\perp}) \simeq \mathbb{C}^N \,,$$

so that  $\mathcal{R}(\tau P_{\lambda_n}) = \mathbb{C}^{2N}_{||}, \ (\mathcal{R}(\tau P_{\lambda_n}))^{\perp} = \mathbb{C}^{2N}_{||}, \ \mathbb{C}^{2N} = \mathbb{C}^{2N}_{||} \oplus \mathbb{C}^{2N}_{\perp}$  and for any linear operator  $L: \mathbb{C}^{2N} \to \mathbb{C}^{2N}$  we can consider the block decomposition  $L = \begin{pmatrix} L_{||} & L_{||\perp} \\ (L_{||\perp})^* & L_{\perp} \end{pmatrix}$ . By using such decompositions in equation (2.6) one gets that  $\mathcal{N}_{\lambda_n}^{\Theta} \neq \{0\}, \ \Theta = B + \bigoplus_{k=1}^{N} P_0$ , if and only if there exists  $\zeta \neq \{0\}, \ \zeta = \zeta_{||} \oplus \zeta_{\perp} \in \mathbb{C}^{2N}_{||} \oplus \mathbb{C}^{2N}_{\perp}$  solving

$$\begin{cases} \zeta_{||} = (B + \bigoplus_{k=1}^{N} P_0 + \lambda_n G_0^* G_{\lambda_n}^{\perp})_{||\perp} \zeta_{\perp}, \\ 0 = (B + \bigoplus_{k=1}^{N} P_0 + \lambda_n G_0^* G_{\lambda_n}^{\perp})_{\perp} \zeta_{\perp}. \end{cases}$$

By (3.2) one obtains  $(\bigoplus_{k=1}^N P_0 + \lambda_n G_0^* G_\lambda^{\perp})_{\perp} = 0$ . Therefore one gets

$$\lambda_n \in \sigma(A^B) \iff \det(B_\perp) = 0.$$

2)  $\Pi \neq 1$ . Given  $\lambda_n \in \mathcal{D}(A)$  we pose

$$\mathcal{R}(\Pi)_{||} := \mathcal{R}(\Pi) \cap (\bigoplus_{k=1}^{N} \operatorname{span}(\hat{\xi}_{n})), \quad \mathcal{R}(\Pi)_{\perp} := \mathcal{R}(\Pi) \cap (\bigoplus_{k=1}^{N} \operatorname{span}(\hat{\xi}_{n}^{\perp})),$$

so that  $\mathcal{R}(\Pi) \cap \mathcal{R}(\tau P_{\lambda_n}) = \mathcal{R}(\Pi)_{||}$ ,  $\mathcal{R}(\Pi) \cap (\mathcal{R}(\tau P_{\lambda_n}))^{\perp} = \mathcal{R}(\Pi)_{\perp}$ ,  $\mathcal{R}(\Pi) = \mathcal{R}(\Pi)_{||} \oplus \mathcal{R}(\Pi)_{\perp}$  and for any linear operator  $L : \mathcal{R}(\Pi) \to \mathcal{R}(\Pi)$  we can consider the block decomposition. position  $L = \begin{pmatrix} \hat{L}_{||} & L_{||\perp} \\ (L_{||\perp})^* & L_{\perp} \end{pmatrix}$ . Define  $\hat{\xi}_{k,n} = \bigoplus_{i=1}^n \hat{\xi}_{i,k,n} \in \mathbb{C}^{2N}$  and  $\hat{\xi}_{k,n}^{\perp} = \bigoplus_{i=1}^n \hat{\xi}_{i,k,n}^{\perp} \in \mathbb{C}^{2N}$ ,  $k = 1, \dots, N$ , by

$$\hat{\xi}_{i,k,n} := \begin{cases} 0, & i \neq k, \\ \hat{\xi}_n, & i = k, \end{cases} \quad \hat{\xi}_{i,k,n}^{\perp} := \begin{cases} 0, & i \neq k, \\ \hat{\xi}_n^{\perp}, & i = k. \end{cases}$$

If  $\Pi \hat{\xi}_{k,n}^{\perp} = 0$  for all k then  $\mathcal{R}(\Pi)_{\perp} = \{0\}$  and in this case

$$\lambda \in \sigma_p(A^{\Pi,B}) \iff \exists k \text{ s.t } \Pi \hat{\xi}_{k,n} = 0.$$

If there exists k' such that  $\Pi \hat{\xi}_{k',n}^{\perp} \neq 0$  then  $\mathcal{R}(\Pi)_{\perp} \neq \{0\}$ . By Remark 2.3

$$\exists k \text{ s.t } \Pi \hat{\xi}_{k,n} = 0 \implies \lambda \in \sigma_p(A^{\Pi,B}).$$

Suppose now  $\Pi \hat{\xi}_{k,n} \neq 0$  for all k, i.e.  $\mathcal{K}_{\lambda_n}^{\Pi} = \{0\}$ . Then, using the above decompositions in equation (2.6) one gets that  $\mathcal{N}_{\lambda_n}^{\Pi,\Theta} \neq \{0\}$ ,  $\Theta = B + \Pi(\bigoplus_{k=1}^N P_0)\Pi$ , if and only if there exists  $\zeta \neq 0$ ,  $\zeta = \zeta_{||} \oplus \zeta_{\perp} \in \mathcal{R}(\Pi)_{||} \oplus \mathcal{R}(\Pi)_{\perp}$  solving

$$\begin{cases} \zeta_{||} = (B + \Pi(\bigoplus_{k=1}^N P_0 + \lambda_n G_0^* G_{\lambda_n}^{\perp}) \Pi)_{||\perp} \zeta_{\perp}, \\ 0 = (B + \Pi(\bigoplus_{k=1}^N P_0 + \lambda_n G_0^* G_{\lambda_n}^{\perp}) \Pi)_{\perp} \zeta_{\perp}. \end{cases}$$

By (3.2) one obtains  $(\Pi(\oplus_{k=1}^N P_0 + \lambda_n G_0^* G_{\lambda_n}^{\perp})\Pi)_{\perp} = 0$ . Therefore one gets, in case there exists k' such that  $\Pi \hat{\xi}_{k',n}^{\perp} \neq 0$  and  $\Pi \hat{\xi}_{k,n} \neq 0$  for all k

$$\lambda_n \in \sigma(A^{\Pi,B}) \iff \det(B_\perp) = 0.$$

#### References

- 1. J. Brüning, V. Geyler, K. Pankrashkin, Spectra of self-adjoint extensions and applications to solvable Schrödinger operators, Rev. Math. Phys. 20 (2008), 1–70.
- 2. Y. Colin De Verdiere, Pseudo-Laplaciens I, Ann. Inst. Fourier 32 (1982), 275–286.
- 3. V. A. Derkach, M. M. Malamud, Generalized resolvents and the boundary value problem for Hermitian operators with gaps, J. Funct. Anal. 95 (1991), 1–95.
- V. Koshmanenko, Singular operators as a parameter of self-adjoint extensions, Oper. Theory Adv. Appl. 118 (2000), 205–223.
- P. Kuchment, Quantum graphs: I. Some basic structures, Waves in Random Media 14 (2004), S107–S128.
- V.A. Mikhailets, A.V. Sobolev, Common eigenvalue problem and periodic Schrödinger operators, J. Funct. Anal. 165 (1999), 150–172.
- A. Posilicano, A Kreĭn-like formula for singular perturbations of self-adjoint operators and applications, J. Funct. Anal. 183 (2001), 109–147.
- 8. A. Posilicano, Boundary triples and Weyl functions for singular perturbations of self-adjoint operators, Methods Funct. Anal. Topology 10 (2004), no. 2, 57–63.
- 9. A. Posilicano, Self-adjoint extensions of restrictions, Operators and Matrices 2 (2008), 483-506.
- 10. P. Šeba, Wave chaos in quantum billiard, Phys. Rev. Lett. 64 (1990), 1855–1858.

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